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## ACTION OF FORCE – FORMALITY OR ESSENCE

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**Abstract.** *It is proved that the author's term "action of force" is not mathematical formality, but it has essential physic significance. This is an expression with whom the action of force can be determined in time. The relation between force and action is established, similar as between force and power. Important examples in analytical, celestial and quantum mechanics are mentioned.*

**1. Introduction.** The concept of *Action* is not uniformly and unambiguously determined in the classical mechanics. Therefore, for the sake of a clearer understanding of this paper, it is necessary to clarify this statement and then to proceed to generalizing and defining the concept of *Action*.

The first concept of the *action* can be found in the work by Leibnitz (1669) [3; pp. 782 - 789], as the *actio formalis*, whose dimension is product of mass, velocity and path. In forth Newton's definition, (1687), it is written that: "Vis impressa est *actio* in corpus exercita, ad mutandum eius statum vel movendi uniformiter in directum" ("An impress'd force is an *actio* exerted upon a body, in order to change its state, either of rest, or of moving uniformly forward in a right line").<sup>1</sup> Accordingly, the force is defined by means of the concept of *action* that has not been previously defined and thus it is assumed to be clear and known. The concept of *actio* was later used by Christian Wolff (1726). Wolff wrote: "*actio* consist of

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<sup>1</sup>Newton Is., *Mathematical Principles of natural Philosophy* (tren. into English by Robert Thorp, MA, London, 1969.

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mass, velocity and space" [3]. Maupertuis P. L. wrote (1746): "The quantity of action is a product of the bodies' masses, their velocities and the distance they are traversing" [3; pp. 53, 881]. Euler [3; pp. 70-75, 791] found (1748) that the sum of all *momentary actions* has the form

$$\int dt \left( \int V dv + V' dv' + V'' dv'' + \dots \right), \quad (*)$$

where  $V, V', V''$  are forces expressed as functions of distances  $v, v', v''$ , ([3], pp. 76, 882). Let's also quote Hamilton who wrote (1834)" that integral

$$A_1 = \int \sum (X' dx + Y' dy + Z' dz) = \int 2E_k dt$$

represents collected living force (vis viva) often termed as *action of system* from its initial to its final position" [3]. In the Van Nostrand's Scientific Encyclopedia, Second Edition (New York, 1947) it is formulated in the following way: "ACTION. In certain discussion of dynamics there is need of an expression for the product of twice the mean total kinetic energy of a system of particles, during a specified interval of time, by the duration of the interval. This product is called *action*. Mathematically, it is expressed by

$$A_1 = 2 \int_{t_0}^{t_1} E_k dt \quad (1)$$

in which  $E_k$  is kinetic energy and  $t_0, t_1$  are the times of beginning and ending of the interval". The significance of formula (1) is also stressed by the fact that it is used to formulate principle of least action that has been formulated and elaborated by the most distinguished and deserving theorists of analytical Mechanics. K. Jacobi, [3; p. 797], even wrote that the principle of least action is the "mother" of the entire analytical mechanics.

In further development of Analytical mechanics the concept of *action* was accepted in the form of the functional

$$A_2 = \int_{t_0}^{t_1} L dt \quad (2)$$

where  $L(q, \dot{q}, t)$  is kinetic potential, often called Lagrange's function.

Physical dimension of *action* is  $ML^2T^{-1}$  (dim mass =  $M$ , dim length =  $L$ , dim time =  $T$ ).

For our purpose here let us also quote a few sentences from A. Sommerfeld's book [4; pp.255,328] " ... just as power is defined as *energy magnitude : time*, so *action* is defined as *energy × time*". ... "When we speak about cause and action, we imply that action is a consequence or result". ... "However, afterwards, if the term action is sanctioned by Helmholtz and Planck, each attempt to substitute it by

some other term will be without perspective”. . . . “As an example, *the elementary quantum of Planck’s action* can be used”.

$$\int pdq = h = 6,6424 \cdot 10^{-34} Js. \quad (3)$$

where  $p$  is generalized impulse and  $q$  is generalized independent coordinate. Let’s quote M. Planck. “Characteristically, the principle of the least action, being fully adapted in the mechanics, by Lagrange, didn’t have any focal and practical influence in the science. It has been observed almost as the mathematical curiosity, as something interesting but, superficial with respect to the Newton’s law of motion. Already in 1837, Poisson called it “only useless rule”. It was not before Thomson and Tait, G.Kirchhoff, K. Neumann, L. Boltzman and other’s who showed it very useful tool in solving several problems in hydrodynamics and elasticity, that the principle began to be appreciable”<sup>2</sup>. Discussing the “element of energy  $E = a\nu$  of the oscillator”, where  $a$  is constant, independent of the nature of oscillator, Planck wrote: “I marked the constant  $a'$  by the letter  $h$  and called it, because it has the dimension of energy and time product, the *elementary quantum of action* or element of action, with difference to the element of energy  $h\nu$ ”. Using measured meaning of constants  $a$  and  $C$ , ( $a' = aC$ ), Planck obtained  $h = 6,55 \cdot 10^{-27} erg \cdot sec$ .

According to the afore mentioned the following terms can be underlined: Leibnitz’s “*actio formalis*”, Newton’s words “*Vis impressa est actio...*”; Maupertuius’ “*quantite d’action*”, Euler’s “*action des forces*”; “*action of systems*” (1), “*Hamilton’s action*” (2); Planck’s “*elementary quantum of action*”. Perhaps this is the reason for the term of “*action*” to be less present in physics, then term “*power*”, which is uniquely defined. There are some other reasons as well. In paper [6] it is shown that functional (2) is not invariant with respect to Lagrange’s generalized coordinates  $q = (q^1, \dots, q^n)^T \in M^n$  and Hamilton’s  $p, q \in T^*M^n$  canonical variables for rheonomic systems. Namely, to put it more clearly, for mechanical systems with time-independent constraints, functional (1) and (2) can be written with respect to Descartes’ coordinates  $y$ , curvilinear coordinates  $x$ , generalized independent coordinates  $q$  as well as with respect to Hamilton’s coordinates  $p, q$  in invariant form

$$A_1 = \int_{t_0}^{t_1} 2E_k dt = \int_{t_0}^{t_1} p_i dq^i = \int_{t_0}^{t_1} a_{ij} \dot{q}^i dq^j = \int_{t_0}^{t_1} p_i \dot{q}^i dt. \quad (4)$$

$$A_2 = \int_{t_0}^t L dt = \int_{t_0}^{t_1} (E_k - E_p) dt = \int_{t_0}^{t_1} (p_i dq^i - H dt); \quad H := E_k + E_p. \quad (5)$$

<sup>2</sup>M.Planck, *Das Princip der kleinsten Wirkung*, “*Die Kultur der Gegenwart*”, T.1, Physik, 195

However, if the constraints of the system depend on time, relation (4) and (5) lose both mathematical and physical sense, since it is

$$A_1 = \int_{t_0}^{t_1} 2E_k dt \neq \int_{t_0}^{t_1} p_i dq^i \neq \int_{t_0}^{t_1} a_{ij} \dot{q}^j dq^i \neq \int_{t_0}^{t_1} p_i \dot{q}^i dt.$$

$$A_2 = \int_{t_0}^{t_1} L dt \neq \int_{t_0}^{t_1} (p_i \dot{q}^i - H) dt, \quad H \neq E_k + E_p.$$

For the above, it was necessary to introduce a general formula of action, which will keep positive mathematical and physical characteristics of defined formulas (1), (2) and (3).

## 2. Definition of the term Action of force.

Here we will retain the accepted dimensions of action, but we will uniformly define the term *action of force*.

**Definition 1.** The Action  $A(\mathbf{F})$  of a force  $\mathbf{F}$  during an interval of time  $t_1 - t_0$  is determined by integral

$$A(\mathbf{F}) := \int_{t_0}^{t_1} W(\mathbf{F}) dt = \int_{t_0}^{t_1} \left( \int_S \mathbf{F} dr \right) dt \quad (6)$$

where  $W(\mathbf{F})$  is the work of all resulting forces  $\mathbf{F}_\nu$  on the  $\nu$ -th material or dynamical point  $M_\nu$  on the path  $S_\nu$  ( $M_\nu \in S_\nu$ );  $ds_\nu = dr_\nu$ . As it is well known (see, for example: [6]) the work of the forces, in general case, is curvilinear integral

$$W(\mathbf{F}) := \int_S \mathbf{F} ds = \int_{S_\nu} \sum_{\nu=1}^N \mathbf{F}_\nu dr_\nu.$$

This definition - just like Euler's action of force (\*) - has more important both mathematical and physical differences with respect to action (2) or (1). The action (6) is invariant with respect to all the above-mentioned transformations, that is

$$A = \int_{t_0}^t W(\mathbf{F}) dt = \int_{t_0}^t W(Y) dt = \int_{t_0}^t W(X) dt = \int_{t_0}^t W(Q) dt =$$

$$\int_{t_0}^t \left( \int_S \mathbf{F} dr \right) dt = \int_{t_0}^t \left( \int_S Y dy \right) dt = \dots = \int_{t_0}^t \left( \int_S Q dq \right) dt$$

both for scleronomic and rheonomic mechanical systems.

For material point of constant mass  $m$ , moving with the velocity  $\mathbf{v}$ ;  $v \in [0, v]$ , it will be

$$A(\mathbf{I}) = \int_{t_0}^{t_1} \left( \int_S -m \frac{d\mathbf{v}}{dt} d\mathbf{r} \right) dt = - \int_{t_0}^{t_1} E_k dt. \quad (7)$$

Although integrals (1) and (2) seem to be similar in form to the integral (7), there is a basic difference between them. Namely, for a constant velocity motion ( $v = \text{const.}$ ) integral (7) is always zero,

$$A(\mathbf{I}) = \int_{t_0}^{t_1} \left( \int_{\mathbf{r}_0}^{\mathbf{r}_1} -m \dot{\mathbf{v}} d\mathbf{r} \right) dt = -m \int_{t_0}^{t_1} \left( \int_{\mathbf{v}_0}^{\mathbf{v}_1} \mathbf{v} d\mathbf{v} \right) dt = - \int_{t_0}^{t_1} (E_k(v_1) - E_k(v_0)) dt = 0,$$

while integral (1) and (2) may attain arbitrarily large values. In additional, integral (1) is the action by definition, and in relation (7) it appears as the consequence of action of inertia force (7).

A very simple **example 1** of a weighty point motion of mass  $m$  at constant velocity  $v = c$  along an ideally smooth horizontal straight line very clearly points out to a quantitative difference between action (1) and action of forces (6). Actions (1) and (2), in this case are

$$A_1 = 2 \int_{t_0}^{t_1} \left( \frac{1}{2} m v^2 \right) dt = m c^2 (t_1 - t_0), \quad (8)$$

$$A_2 = \int_{t_0}^{t_1} L dt = m c^2 \frac{t_1 - t_0}{2}. \quad (9)$$

However, action (7) is

$$A(\mathbf{I}) = \int_{t_0}^{t_1} \left( -m \frac{d\mathbf{v}}{dt} \right) dt = 0, \quad (10)$$

since in this example  $\frac{dv}{dt} = 0$ . It is enough to give some thought to result (9) in order to reach the conclusion that such an action does not exist in nature while it is senseless in theory. Classic terms of action (1) and (2), differ from the new introduced ones, (6) and (7), not only by quantity, but also logically, which is seen from the following text. Action of a force should be taken into account in considering the motion of a mechanical system. If the trajectory is given by parametrized finite equations of the form  $\mathbf{r} = \mathbf{r}(\lambda, t)$  or by differential equations of motion  $m\dot{\mathbf{v}} = \mathbf{F}$ , then the curvilinear integral of work  $W(\mathbf{F})$  can be reduced to a definite integral

$$W = \int_S \mathbf{F}(\mathbf{r}, \mathbf{v}, t) d\mathbf{r} = \int_{t_0}^t \mathbf{F}(\mathbf{r}(t), \mathbf{v}(t), t) \mathbf{v}(t) dt = \int_{t_0}^t \mathcal{F}(t) dt.$$

That corresponds to known relations

$$\frac{dW}{dt} = \mathcal{F} = \mathbf{F} \frac{d\mathbf{r}}{dt} \longrightarrow dW = \mathbf{F} d\mathbf{r}.$$

**Example 2.** Material point of the mass  $m$  moves along the spiral:  $x = a \cos \omega t, y = b \sin \omega t, z = ct$ . Let us find action of the force in the interval  $[t_0 = 0, t_1 = 2\pi]$ . Since  $\ddot{x} = -m\omega^2 a \cos \omega t = -m\omega^2 x; \ddot{y} = -m\omega^2 b \sin \omega t = -m\omega^2 y; \ddot{z} = 0$ , we have  $F_x = -m\omega^2 x, F_y = -m\omega^2 y, F_z = 0$ . So:

$$W(\mathbf{F}) = \int_S \mathbf{F} d\mathbf{r} = -m\omega^2 \left( a^2 \int_0^{2\pi} \cos \omega t d \cos \omega t + b^2 \int_0^{2\pi} \sin \omega t d \sin \omega t \right) = 0;$$

and then

$$A(\mathbf{F}) = 0.$$

**Example 3.** The force  $\mathbf{F} = -\mu \mathbf{v}$  acting on a material point of masse  $m$ . Action of force  $\mathbf{F}$  is

$$A(\mathbf{F}) = \int_{t_0}^{t_1} \left( \int_S (-\mu \mathbf{v} d\mathbf{r}) \right) dt.$$

Since from the differential equation of motion  $m\dot{\mathbf{v}} = -\mu \mathbf{v}$  follows that  $\mathbf{v} = \mathbf{v}_0 - \frac{\mu}{m}(\mathbf{r} - \mathbf{r}_0)$ , we get

$$A(\mathbf{F}) = - \int_{t_0}^t \left( \mu \int_{\mathbf{r}_0}^{\mathbf{r}} \left[ \mathbf{v}_0 - \frac{\mu}{m}(\mathbf{r} - \mathbf{r}_0) \right] d\mathbf{r} \right) dt. \quad (11)$$

### 3. Planch's elementary quantum of action

Although, Quantum of Physics is based on the standpoint (see for example [9]) that using the classical Mechanics it is not possible to explain the base of Theory of quantum, "The hypothesis of elementary quantum of action is a good step in explaining physical phenomena, as it assumes finite domains, instead of infinitesimal once:

$$\int \int dq dp = h. \quad (12)$$

where  $q$  is a generalized independent coordinate of Physical systems and  $p$  corresponding impuls (momentum); quantity  $h$  is *the elementary quantum of action* and presented by universal constant with dimension of energy  $\times$  time" (M. Planch, [2; pp. 286-287]). Relation (12), using the Green's theorem can be written as:

$$A = \int \int dq dp = \int pdq = \int_{t_0}^{t_1} 2E_k dt = h. \quad (13)$$

In the further text on model of atom and physical data let's us Berkely Physics course – Quantum Physics (IV) by Wikhmann [9; Sec. 2.29]. "Our discussion is therefore very defective, but it nevertheless gives us some kind of picture of the nature of heavy atoms. It follows from this picture that transitions in the state of motion of the outermost or optical electrons will involve energies of the order of an electron volt, which roughly corresponds to wavelengths of the emitted photons in the optical region,...". Electron mass  $m_e = (9.10908 \pm 0,00013) \cdot 10^{-31} \text{kg}$ ; velocity of light  $c = (997925 \pm 0,00001) \cdot 10^8 \text{ms}^{-1}$ . With numerical values of mass, velocity and time, or impulse and wavelength, Planck's elementary quantum of action can be calculated. I will presume that the carrier of action is the part of electron of mass  $\Delta m = 1.3 \cdot 10^{-35} \text{kg}$ , as error of measuring of the electron mass. Let's call it a "crumb". At direct calculation of the quantum of action with formula (1) and (2) the problem of limits of the interval  $t_1 - t_0$  appears. However, the integral of action of force [6] removes that problem, as  $W$  is a work of force which is expelling a "crumb" from atom. This is in correspondence with Einstein's explanation of photoeffect [10; Sec.42]. If a particle expels from a body, with velocity  $u$  as it is well known, the reactive force  $F = \dot{m}u$  is appearing. At the expelling crumb of mass  $\Delta m$ , the reactive force  $\dot{m}u$  is acting on particle, while the force  $-\dot{m}u$  is acting on atom, in correspondence of the principle of acting and counteraction. The time of duration of action of force is the time necessary for a crumb to accelerate from zero velocity to the velocity  $c$ . If, as example, we assume that the wavelength of mass emission  $\Delta m = 0,00013 \cdot 10^{-31} \text{kg}$  is equal to  $\lambda = 1,7 \cdot 10^{-7} \text{m}$ , then the interval of time of crumb expelling, on light speed  $c = 2,997926 \cdot 10^8 \text{ms}^{-1}$  is equal

$$\tau = \frac{\lambda}{c} = 0,567 \cdot 10^{-15} = 5,67 \cdot 10^{-16} \text{s}.$$

Thus, the action of this crumb is

$$\begin{aligned} A &= \int_0^{\tau} \left( \int_m^{m-\Delta m} u^2 d(\Delta m) \right) dt = \int_0^{\tau} (\Delta m u^2) dt = \Delta m u^2 \tau \\ &= 1,3 \cdot 10^{-35} \cdot 2,997926^2 \cdot 10^{16} \cdot 5,67 \cdot 10^{-16} = 6.6247 \cdot 10^{-34} \text{Js}, \end{aligned} \quad (14)$$

which corresponds to the Planck's elementary action.

**4 De Broglie's relation** At the similar but more general level, de Broglie's relation [9; Sec 5.23, p 196]

$$h = \lambda p, \quad (15)$$

where  $p$  is the momentum of the particle, and  $\lambda$  its de Broglie wave length can be derived. Let us show this again on the classical way. If a particle of mass  $m$  makes a road of  $c$  meters for 1 second, then the wave of wavelength  $\lambda$  will make it for time  $\tau = \frac{\lambda}{c}$ . With substitution in the relation (3), relation

$$h = \int_0^{\tau} 2E_k dt = mc^2 \tau = mc^2 \frac{\lambda}{c} = p\lambda \quad (16)$$

is obtained, which is de Broglie's relation (15).

"Both  $p$  and  $\lambda$  are independently measurable quantities, and by measuring a pair of corresponding variables  $(p, \lambda)$  we can determine Planck's constant  $h$ . It is remarkable empirical fact that we always get the same value for  $h$ , irrespective of which kind of particle we are observing, and that this is so is not a triviality" [10, p. 200].

### 5. Classical mechanical approach to the basic principle of quantum physics

Let's quote the first Berkeley Physics course—Quantum Physics ([9] Sec. 1.46). "As we will learn later the relation (46a;  $\frac{E}{\nu} = X_2$ ) expresses a very basic principle of quantum physics, namely that energy and frequency are *universally* related by  $E = h\nu$ . Such a relationship is entirely foreign to classical physics, and the mysterious constant  $X_2 (= h)$  in (46a) is a manifestation of secrets of nature unsuspected at the time."

Solving the problem in general form with canonical variables,  $p := (p_1, \dots, p_n)$ ,  $q := (q^1, \dots, q^n)$  Hamilton wrote the action (2) in the form:

$$A_2 = \int_{t_0}^{t_1} L(q, \dot{q}) dt = \int_{t_0}^{t_1} [2E_k(q, p) - E(p, q)] dt. \quad (17)$$

or, regarding relation (1) or (13) and (6),

$$A_2 = A_1 - \int_0^\tau E dt. \quad (18)$$

For conservative system, which kinetical potential  $L = E_k - E_p = 0$ , it follows

$$A_2 = h - E\tau = 0, \quad (19)$$

or

$$E = h\tau^{-1}. \quad (20)$$

For an oscillatory process, in which the frequency  $\nu = \tau^{-1}$ , the basic relation of quantum physics is:

$$E = h\nu. \quad (21)$$

#### Example 4. Harmonic oscillator.

Let's define the action (3) for harmonic oscillator.

$$\ddot{x} + \omega^2 x = 0, \omega^2 = \frac{c}{m}, c > 0, m > 0, x(t_0) = x_0, \dot{x}(0) = \dot{x}_0.$$



Finally equation of motion is

$$x(t) = x_0 \cos(\omega t) + \frac{\dot{x}}{\omega} \sin(\omega t).$$

As the  $p(t) = m\dot{x}$ , introducing it to the relation (3) it follows

$$h = \int_0^\tau m\dot{x}dx = \int_0^\tau m\dot{x}^2 dt = \int_0^\tau \frac{m}{2}(\dot{x}_0^2 + \omega x_0^2)dt = \int_0^\tau (E_k + E_p)dt = E\tau, \quad (22)$$

because

$$E_k = \frac{m}{2}\dot{x}^2, \quad E_p = \frac{c}{2}x^2; \quad E_k + E_p = E = \text{const.}$$

Further, as

$$\tau = \nu^{-1} = \frac{2\pi}{\omega}$$

we obtain:

$$E = h\nu = \hbar\omega$$

where:

$$\hbar = \frac{h}{2\pi} = 1.05449 \cdot 10^{-34} \text{ Js.}$$

We can observe that from the previous equality it follows

$$\int_0^\tau m\dot{x}dx - (E_k + E_p)dt = \int_0^\tau Ldt = 0.$$

The qualitative analysis of this relation shows that separate phase trajectories correspond to the zero kinetical potential  $L = 0$ . This also corresponds to the crumb expelling from the atom after the cyclus  $\nu$  is finished.

## 6. Two bodies model

In scientific literatur it is accepted that the gravitation force between two bodies of the mass  $m_1$  and  $m_2$  is defined by formula

$$F = -\kappa \frac{m_1 m_2}{\rho^2} \quad (23)$$

Using the variational Maupertuis–Lagrange’s principle Ji-Huan He obtained a more general result<sup>3</sup>

$$F = -\chi \frac{m_1 m_2}{\rho}. \quad (24)$$

<sup>3</sup>See his manuscript "A variational approach to problem of two bodies"

Let's check this result by using our general principle of action on the problem of two bodies. In the paper [8] and monography [6], the principle of action is formulated by the relation:

$$\delta \int_{t_0}^{t_1} [W(\mathbf{F}) - W(\mathbf{I})] dt = 0 \quad (25)$$

where  $W(\mathbf{F})$  is work of force  $\mathbf{F}$  and  $W(\mathbf{I})$  is work of inertia force  $\mathbf{I}$ . Distance  $\rho(t)$  between bodies is changes in time and thus it can be written as

$$f_1 = \left[ (x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2 \right]^{\frac{1}{2}} - \rho(t) = 0. \quad (26)$$

Thus it is only possible to claim, that beside geometrical equation (26), forces of inertia also exist. As the negative work of inertia force is equal to kinetic energy [6], relation (25) by using the boundary conditions  $\delta x(t_0) = \delta x(t_1) = \dots \delta z(t_1) = 0$ , is reduced to:

$$- \int_{t_0}^{t_1} [m_1(\ddot{x}_1 \delta x_1 + \ddot{y}_1 \delta y_1 + \ddot{z}_1 \delta z_1) + m_2(\ddot{x}_2 \delta x_2 + \dots + \ddot{z}_2 \delta z_2)] dt = 0. \quad (27)$$

With this, the equation

$$\delta f_1 = (x_2 - x_1) \delta(x_2 - x_1) + \dots + (z_2 - z_1) \delta(z_2 - z_1) = 0. \quad (28)$$

should be satisfied. If the equation is integrated in the interval  $(t_1 - t_0)$ , multiplied by Lagrange's multiplier  $\lambda$ , and then summed with equation (27), it is obtained

$$\lambda_1 \int_{t_0}^{t_1} [(x_2 - x_1) \delta(x_2 - x_1) + \dots + (z_2 - z_1) \delta(z_2 - z_1)] dt - \int_{t_0}^{t_1} [m_1 \ddot{x}_1 \delta x_1 + \dots + m_2 \ddot{x}_2 \delta x_2] dt = 0.$$

This variational relation is equivalent to the following system of differential equations of motion of two bodies

$$\begin{aligned} m_1 \ddot{x}_1 &= -\lambda(x_2 - x_1), \\ m_1 \ddot{y}_1 &= -\lambda(y_2 - y_1), \\ m_1 \ddot{z}_1 &= -\lambda(z_2 - z_1); \\ m_2 \ddot{x}_2 &= \lambda(x_2 - x_1), \\ m_2 \ddot{y}_2 &= \lambda(y_2 - y_1), \\ m_2 \ddot{z}_2 &= \lambda(z_2 - z_1). \end{aligned}$$

Using these equations and conditions of acceleration the multiplier  $\lambda$  can be defined as:

$$\lambda = \chi \frac{m_1 m_2}{\rho^2}, \quad ()$$

where:

$$\chi = \frac{\dot{\rho}^2 + \rho \ddot{\rho} - v_{or}^2}{m_1 + m_2},$$

and

$$v_{or}^2 := (\dot{x}_2 - \dot{x}_1)^2 + (\dot{y}_2 - \dot{y}_1)^2 + (\dot{z}_2 - \dot{z}_1)^2.$$

Thus, the gravitation force of the two bodies of mass  $m_1$  and  $m_2$  is

$$F = \chi \frac{m_1 m_2}{\rho}.$$

The same result is obtained with precondition that the motion of body has to satisfy the second Kepler's law, we can write in the form

$$[\rho, v_{or}] = \begin{vmatrix} \mathbf{e}_1 & \mathbf{e}_2 & \mathbf{e}_3 \\ x_2 - x_1 & y_2 - y_1 & z_2 - z_1 \\ \dot{x}_2 - \dot{x}_1 & \dot{y}_2 - \dot{y}_1 & \dot{z}_2 - \dot{z}_1 \end{vmatrix} = \mathbf{c} = c_1 \mathbf{e}_1 + c_2 \mathbf{e}_2 + c_3 \mathbf{e}_3 = 0.$$

That means that beside equations (26) there are additional three equations

$$\begin{aligned} f_2 &= (y_2 - y_1)(\dot{z}_2 - \dot{z}_1) - (\dot{y}_2 - \dot{y}_1)(z_2 - z_1) = c_1, \\ f_3 &= (z_2 - z_1)(\dot{x}_2 - \dot{x}_1) - (\dot{z}_2 - \dot{z}_1)(x_2 - x_1) = c_2, \\ f_4 &= (x_2 - x_1)(\dot{z}_2 - \dot{z}_1) - (\dot{z}_2 - \dot{z}_1)(y_2 - y_1) = c_3. \end{aligned}$$

Integrating the equation of variation

$$\delta f_2 = z \delta y + y \delta z - z \delta y - y \delta z = 0,$$

where  $x := x_2 - x_1$ ,  $y := y_2 - y_1$ ,  $z := z_2 - z_1$ , we obtain:

$$\int_{t_0}^{t_1} \delta f_2 dt = 2 \int_{t_0}^{t_1} (z \delta y - y \delta z) dt = 0,$$

and then

$$\int_{t_0}^{t_1} \delta f_3 dt = 2 \int_{t_0}^{t_1} (x \delta z - z \delta x) dt = 0$$

$$\int_{t_0}^{t_1} \delta f_4 dt = 2 \int_{t_0}^{t_1} (y \delta x - x \delta y) dt = 0.$$

With successive multiplying of these relations by multipliers  $\beta_2, \beta_3$  and summing with equation (27), it follows

$$\begin{aligned} & \lambda \int_{t_0}^{t_1} (x\delta x + y\delta y + z\delta z) dt + \beta_1 \int_{t_0}^{t_1} (\dot{x}\delta y - \dot{y}\delta x) dt \\ & + \beta_2 \int_{t_0}^{t_1} (\dot{y}\delta z - \dot{z}\delta y) dt + \beta_3 \int_{t_0}^{t_1} (\dot{z}\delta x - \dot{x}\delta z) dt \\ & - \int_{t_0}^{t_1} [m_1(\ddot{x}_1\delta x_1 + \ddot{y}_1\delta y_1 + \ddot{z}_1\delta z_1) + m_2(\ddot{x}_2\delta x_2 + \ddot{y}_2\delta y_2 + \ddot{z}_2\delta z_2)] dt = 0. \end{aligned}$$

Remembering that  $x := x_2 - x_1, \dots, z := z_2 - z_1$  follows 6 equations of motion follow:

$$\begin{aligned} m_1\ddot{x}_1 &= -\lambda x + \beta_2\dot{z} - \beta_3\dot{y}, \\ m_1\ddot{y}_1 &= -\lambda y - \beta_1\dot{z} + \beta_3\dot{x}, \\ m_1\ddot{z}_1 &= -\lambda z + \beta_1\dot{y} - \beta_2\dot{x}, \\ m_2\ddot{x}_2 &= \lambda x - \beta_2\dot{z} + \beta_3\dot{y}, \\ m_2\ddot{y}_2 &= \lambda y - \beta_3\dot{x} + \beta_1\dot{z}, \\ m_2\ddot{z}_2 &= \lambda z - \beta_1\dot{y} + \beta_2\dot{x}. \end{aligned} \tag{29}$$

Substituting accelerations  $\ddot{x}, \ddot{y}, \ddot{z}$  from equations (29) into the condition of accelerations

$$\begin{aligned} \dot{f}_1 &= x(\ddot{x}_2 - \ddot{x}_1) + y(\ddot{y}_2 - \ddot{y}_1) + z(\ddot{z}_2 - \ddot{z}_1) + v_{or}^2 - \dot{\rho}^2 - \rho\ddot{\rho} = 0, \\ \dot{f}_2 &= y(\ddot{z}_2 - \ddot{z}_1) - z(\ddot{y}_2 - \ddot{y}_1) = 0, \\ \dot{f}_3 &= z(\ddot{x}_2 - \ddot{x}_1) - x(\ddot{z}_2 - \ddot{z}_1) = 0, \\ \dot{f}_4 &= x(\ddot{y}_2 - \ddot{y}_1) - y(\ddot{x}_2 - \ddot{x}_1) = 0, \end{aligned}$$

it is obtained:

$$\begin{aligned} \beta_2 &= \beta_3 = \beta_4 = 0; \\ \lambda &= \frac{\dot{\rho}^2 + \rho\ddot{\rho} - v_{or}^2}{m_1 + m_2} \frac{m_1 m_2}{\rho^2}, \end{aligned}$$

where:

$$v_{or}^2 := \dot{x}^2 + \dot{y}^2 + \dot{z}^2.$$

This corresponds to Vujičić's result, obtained by principle of equilibrium [6],[7].

Finally, let's get back to the title: Action of force – formality or essence. On the example of inertial motion, a difference between action (1) and action of inertia force (7) was clearly observed. In the action (2), the integral

$$\int_{t_0}^{t_1} E_p dt \tag{30}$$

appears, where  $E_p$  is potential energy. In many books (see for example [1]) in mechanics and physics, it is learnt that the gravitational potential is

$$E_p = -G \frac{m_1 m_2}{\rho}$$

Assuming that the satellite of mass  $m_2$  is moving on the circular orbit at the distance  $\rho = R = \text{const.}$  from the centre of inertia of the body of mass  $m_1$  it follows that the action (30) is

$$\int_{t_0}^{t_1} G \frac{m_1 m_2}{R} (t_1 - t_0) = G \frac{m_1 m_2}{R} (t_1 - t_0).$$

As it is seen, it will increase, which has no physical sense. On the other hand, for the same bodies and motion at the gravitation force

$$F = \chi \frac{m_1 m_2}{\rho}$$

the action of force (6) will be equal to zero.

**Conclusion.** It is proved that the author's term "action of force" is not mathematical formality, but it has essential physical significance. This is an expression with which the action of force can be determined in time. The relation between force and action is established, similar as between force and the power. Important examples in analytical, celestial and quantum mechanics are mentioned.

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## DEJSTVO SILE-FORMALNOST ILI SUŠTINA

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Pojam "Dejstvo" nije jedinstveno i jednoznačno definisan u klasičnoj mehanici. Postoje razlike i u matematičkom i u fizičkom smislu između pojedinih definicija dejstva. Zato se ponekad postavlja pitanje o formalnosti ili suštini ovog pojma u analitičkoj mehanici. Reč "Dejstvo" (lat. Actio) u fizici prvo nalazimo kod Leibnitz-a kao "Actio formalis" (1669). kasnije (1687) Is. Newton je, definišući sile napisao: "Vis impressa est actio in corpus exercita",... (Nametnuta ili pokretačka sila je dejstvo...)." O razvoju pojma dejstva dovoljno i iscrpno vidi u antologiji "Varijacioni principi mehaniki" (na ruskom jeziku) (1959) od Polaka. Uočiće se i pojmovi: Mopertuis-ova količina dejstva, Euler-ovo dejstvo sila, "Dejstvo po Lagrange-u", Hamilton-ovo dejstvo,... Planck-ovo elementarno dejstvo. Možda je ta raznovrsnost razlog što pojam dejstva nije usvojen u fizici i mehanici, kao što je to pojam snage.

Ovde je redefinisana pojam "dejstvo sile" sa pretenzijom da postane jedinstven u fizici i mehanici. Kao i u većini dosadašnjih definicija usvojena je fizička dimenzija dejstva  $ML^2T^{-1}$ , a "dejstvo sile" je definisano kao integral

$$\mathbf{A} = \int_{t_0}^t \mathbf{W}(\mathbf{F}) dt = \int_{t_0}^t \mathbf{W}(Y) dt = \int_{t_0}^t \mathbf{W}(x) dt = \int_{t_0}^t \mathbf{W}(\mathbf{Q}) dt = \int_{t_0}^t (\int_S \mathbf{F} d\mathbf{r}) dt = \int_{t_0}^t (\int_S \mathbf{F} d\mathbf{r}) dt = \int_{t_0}^t (\int_S Y dy) dt = \int_{t_0}^t (\int_S Q dq) dt \quad (1)$$

gde su  $x, y, q$  u istom redosledu Descartes-ove koordinate, krivolinijske i generalisane koordinate;  $\mathbf{W}(\mathbf{F})$  je rad sile  $\mathbf{F}$ . Polazeći od definicije dejstva sile i principa dejstva i protivdejstva, određuje se na klasičan način elementarno Planck-ovo dejstvo:

$$\mathbf{A}(\mathbf{F}) = -\mathbf{A}(\mathbf{I}) = h = 6.62 \cdot 10^{-34} \text{ Js} \quad (2)$$

gde je  $h$  Planck-ovo elementarno dejstvo. Za klasični harmonijski oscilator lako se dobija osnovna relacija kvantne teorije:

$$\mathbf{E} = h\nu \quad (3)$$

Sve to kao i neke numeričke vrednosti u navedenim primerima, pokazuje da nova formula dejstva sila nije matematički formalizam, nego ima suštinski značaj, ne manji od pojma snage. Na osnovu te definicije formulisan je generalisani integralni varijacioni princip dejstva relacijom:

$$\delta \int_{t_0}^t [\mathbf{W}(\mathbf{Q}) - \mathbf{W}(\mathbf{I})] dt = 0$$

gde je  $\mathbf{W}(\mathbf{I})$  rad generalisane sile inercije  $\mathbf{I}$ .