# ON GENERALIZED BERTRAND CURVES IN EUCLIDEAN 3-SPACE 

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#### Abstract

In this paper, we generalize the notion of Bertrand curve in Euclidean 3 -space analogously as in Minkowski 3-space. According to this generalization, the Bertrand curve conditions of a given space curve are obtained and the relations between Frenet vectors and curvature functions are revealed. A curve that meets the given condition is constructed as an example. Keywords: Bertrand curve, Euclidean 3-space, Frenet vectors, curvature functions.


## 1. Introduction

In the theory of curves in Euclidean space, one of the important and interesting problems is characterization of a regular curve. In the solution of the problem, the curvature functions $k_{1}$ (or $\varkappa$ ) and $k_{2}$ (or $\tau$ ) of a regular curve have an effective role. For example: if $k_{1}=0=k_{2}$, then the curve is a geodesic or if $k_{1}=$ constant $\neq 0$ and $k_{2}=0$, then the curve is a circle with radius $1 / k_{1}$, etc. Another way in the solution of the problem is the relationship between the Frenet vectors and Frenet planes of the curves $[6,10]$. An interesting example of relations between Frenet vectors belonging to pairs of space curves is Bertrand curves.

A Bertrand curve is a curve in the Euclidean space such that its principal normal is the principal normal of the second curve [2,12]. The study of this kind of curves has been extended to many other ambient spaces. In [9], Pears studied this problem

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for curves in the $n$-dimensional Euclidean space $\mathbb{E}^{n}, n>3$, and showed that a Bertrand curve in $\mathbb{E}^{n}$ must belong to a three-dimensional subspace $\mathbb{E}^{3} \subset \mathbb{E}^{n}$. This result is restated by Matsuda and Yorozu [8]. They proved that there was not any special Bertrand curves in $\mathbb{E}^{n}(n>3)$ and defined a new kind, which is called (1,3)-type Bertrand curves in 4-dimensional Euclidean space. Bertrand curves and their characterizations were studied by many researchers in Minkowski 3 -space and Minkowski space-time $[1,3,5]$ as well as in Euclidean space. In addition, (1, 3)-type Bertrand curves were studied in semi-Euclidean 4 -space with index 2 [11].

A new generalization for Bertrand curves is given by Zhang and Pei in 2020 [13]. In this study, instead of classical condition for Bertrand curves, generalized Bertrand curves are defined such that the principal normal of a given curve belongs to a normal space of another curve.

In this paper, we generalize the notion of Bertrand curve in Euclidean 3-space analogously as in Minkowski 3 -space. According to this generalization, the Bertrand curve conditions of a given space curve are obtained and the relations between Frenet vectors and curvature functions are revealed. An example of a curve is constructed that satisfies the given conditions.

## 2. Preliminaries

In this section, we give some well known results from Euclidean geometry [4, 7]. Let $\mathbb{E}^{3}$ be the 3 -dimensional Euclidean space equipped with the inner product $\langle X, Y\rangle=$ $x_{1} y_{1}+x_{2} y_{2}+x_{3} y_{3}$, where $X=\left(x_{1}, x_{2}, x_{3}\right)$ and $Y=\left(y_{1}, y_{2}, y_{3}\right) \in \mathbb{E}^{3}$. The norm of $X$ is given by $\|X\|=\sqrt{\langle X, X\rangle}$ and the vector product is given by

$$
X \times Y=\left(\begin{array}{lll}
e_{1} & e_{2} & e_{3} \\
x_{1} & x_{2} & x_{3} \\
y_{1} & y_{2} & y_{3}
\end{array}\right)
$$

where $\left\{e_{1}, e_{2}, e_{3}\right\}$ is the canonical basis of $\mathbb{E}^{3}$.
Let $I$ be an interval of $\mathbb{R}$ and let $\gamma: I \rightarrow \mathbb{E}^{3}$ be a regular space curve, that is, $\gamma^{\prime}(t) \neq 0$ for all $t \in I$, where $\gamma^{\prime}(t)=\frac{d \gamma}{d t}(t)$. We say that $\gamma$ is nondegenerate condition if $\gamma^{\prime}(t) \times \gamma^{\prime \prime}(t) \neq 0$ for all $t \in I$. If we take the arc-length parameter $s$, that is, $\left\|\gamma^{\prime}(s)\right\|=1$ for all $s$, then the tangent vector , the principal normal vector , and the binormal vector are given by

$$
\begin{aligned}
T(s) & =\gamma^{\prime}(s) \\
N(s) & =\frac{\gamma^{\prime \prime}(s)}{\left\|\gamma^{\prime \prime}(s)\right\|} \\
B(s) & =T(s) \times N(s)
\end{aligned}
$$

where $\gamma^{\prime}(s)=\frac{d \gamma}{d s}(s)$. Then $\{T(s), N(s), B(s)\}$ is a moving frame of $\gamma(s)$ and we have the Frenet-Serret formula:

$$
\left[\begin{array}{l}
T^{\prime}  \tag{2.1}\\
N^{\prime} \\
B^{\prime}
\end{array}\right]=\left[\begin{array}{ccc}
0 & \kappa(s) & 0 \\
-\kappa(s) & 0 & \tau(s) \\
0 & -\tau(s) & 0
\end{array}\right]\left[\begin{array}{l}
T \\
N \\
B
\end{array}\right]
$$

where

$$
\begin{aligned}
\kappa(s) & =\left\|\gamma^{\prime \prime}(s)\right\| \\
\tau(s) & =\frac{\operatorname{det}\left(\gamma^{\prime}(s), \gamma^{\prime \prime}(s), \gamma^{\prime \prime \prime}(s)\right)}{\kappa^{2}(s)}
\end{aligned}
$$

If we take general parameter $t$, then the tangent vector , the principal normal vector and the binormal vector are given by

$$
\begin{aligned}
T(t) & =\frac{\dot{\dot{\gamma}}(t)}{\|\dot{\dot{\gamma}}(t)\|} \\
N(t) & =B(t) \times T(t) \\
B(t) & =\frac{\dot{\gamma}(t) \times \dot{\gamma}(t)}{\|\dot{\gamma}(t) \times \dot{\gamma}(t)\|}
\end{aligned}
$$

where $\dot{\gamma}(t)=\frac{d \gamma}{d t}(t)$. Then $\{T(t), N(t), B(t)\}$ is a moving frame of $\gamma(t)$ and we have the Frenet-Serret formula:

$$
\left[\begin{array}{c}
\dot{T}(t)  \tag{2.2}\\
\dot{N}(t) \\
\dot{B}(t)
\end{array}\right]=\left[\begin{array}{ccc}
0 & \|\dot{\gamma}(t)\| \kappa(t) & 0 \\
-\|\dot{\gamma}(t)\| \kappa(t) & 0 & \|\dot{\gamma}(t)\| \tau(t) \\
0 & -\|\dot{\gamma}(t)\| \tau(t) & 0
\end{array}\right]\left[\begin{array}{c}
T(t) \\
N(t) \\
B(t)
\end{array}\right]
$$

where

$$
\begin{aligned}
\kappa(t) & =\frac{|\dot{\gamma}(t) \times \ddot{\gamma}(t)|}{|\dot{\gamma}(t)|^{3}} \\
\tau(t) & =\frac{\operatorname{det}(\dot{\gamma}(t), \ddot{\gamma}(t), \dddot{\gamma}(t))}{|\dot{\gamma}(t) \times \ddot{\gamma}(t)|^{2}}
\end{aligned}
$$

## 3. Generalized Bertrand curves in Euclidean 3-space

In this section, generalized Bertrand curve concept given by Zhang and Pein (see [13]) Minkowski 3-space will be defined in 3-dimensional Euclidean space $\mathbb{E}^{3}$ and generalized Bertrand curve conditions will be obtained for a curve in this space.

Definition 3.1. Let $\alpha(s)$ and $\alpha^{*}\left(s^{*}\right)$ be two curves in 3-dimensional Euclidean space. If the principal normal $N(s)$ of $\alpha(s)$ lies in the normal plane of $\alpha^{*}\left(s^{*}\right)$ and the angle between $N(s)$ and $N^{*}\left(s^{*}\right)$ is $\theta$ at the corresponding points, then we call $\alpha(s)$ a generalized Bertrand curve, $\alpha^{*}\left(s^{*}\right)$ is a generalized Bertrand mate of $\alpha(s)$. Also $\left(\alpha(s), \alpha^{*}\left(s^{*}\right)\right)$ is called a pair of generalized Bertrand curves.

According to the above definition if $\alpha(s)$ is a generalized Bertrand Curve in $\mathbb{E}^{3}$ then the following holds :
(i) $\quad N=\cos \theta N^{*}\left(s^{*}\right)+\sin \theta B^{*}\left(s^{*}\right)$,
(ii) $\left\langle N(s), N^{*}\left(s^{*}\right)\right\rangle=\cos \theta=$ constant
where $N^{*}$ and $B^{*}$ are principal normal and binormal vectors of $\alpha^{*}$.
Theorem 3.1. Let $\alpha(s)$ be a generalized Bertrand curve in Euclidean 3-space, parametrized by its arc-length $s$ and $\alpha^{*}\left(s^{*}\right)$ be the generalized Bertrand mate curve of $\alpha(s)$ in $\mathbb{E}^{3}$ such that the principal normal $N(s)$ of $\alpha(s)$ lies in the normal plane spanned by $\left\{N^{*}, B^{*}\right\}$ and the angle between $N$ and $N^{*}$ is $\theta$ at the corresponding points. The curvatures and Frenet vector of $\alpha$ and $\alpha^{*}$ are related as follows:

$$
\begin{aligned}
T^{*} & =\left(\frac{1-\lambda \kappa}{f^{\prime}}\right) T+\frac{\lambda \tau}{f^{\prime}} B, \\
N^{*} & =e T+\cos \theta N+h B \\
B^{*} & =\frac{D_{1}}{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}} T+\frac{D_{2}}{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}} N+\frac{D_{3}}{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}} B
\end{aligned}
$$

and

$$
\kappa^{*}\left(s^{*}\right)=\frac{\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)}{\left[(1-\lambda \kappa)^{2}+\lambda^{2} \tau^{2}\right] \cos \theta}, \quad \tau^{*}\left(s^{*}\right)=\frac{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}}{f^{\prime}}
$$

where $\cos \theta \neq 0$ and $\lambda \in \mathbb{R}_{0},\left(f^{\prime}\right)^{2}=(1-\lambda \kappa)^{2}+\lambda^{2} \tau^{2}$,

$$
e=\frac{-\cos \theta\left(\lambda \kappa^{\prime}+\frac{f^{\prime \prime}}{f^{\prime}}(1-\lambda \kappa)\right)}{\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)} \quad, \quad h=\frac{\cos \theta\left(\lambda \tau^{\prime}+\frac{f^{\prime \prime}}{f^{\prime}} \lambda \tau\right)}{\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)}
$$

and

$$
D_{1}=e^{\prime}-\kappa \cos \theta+\kappa^{*}(1-\lambda \kappa), \quad D_{2}=e \kappa-h \tau, \quad D_{3}=\tau \cos \theta+h^{\prime}+\kappa^{*} \lambda \tau
$$

Proof. Assume that there exists the generalized Bertrand curve $\alpha$ in $\mathbb{E}^{3}$ and its $\alpha^{*}$ generalized Bertrand mate $\alpha^{*}$ in $\mathbb{E}^{3}$. Then $\alpha^{*}$ can be parametrized by

$$
\begin{equation*}
\alpha^{*}\left(s^{*}\right)=\alpha(s)+\lambda(s) N(s) \tag{3.1}
\end{equation*}
$$

where $s^{*}=s^{*}(s)$. Differentiating equation (3.1) with respect to $s$ and using Frenet frame (2.1), we get

$$
T^{*} f^{\prime}=(1-\lambda \kappa) T+\lambda^{\prime} N+\lambda \tau B
$$

By taking the inner product of the last relation by $N=\cos \theta N^{*}+\sin \theta B^{*}$, we have $\lambda^{\prime}=0$.Substituting this in the last relation, we find

$$
\begin{equation*}
T^{*} f^{\prime}=(1-\lambda \kappa) T-\lambda \tau B \tag{3.2}
\end{equation*}
$$

From equation (3.2), we obtain

$$
\begin{equation*}
\left\langle T^{*} f^{\prime}, T^{*} f^{\prime}\right\rangle=\left(f^{\prime}\right)^{2}=(1+\lambda \kappa)^{2}+(\lambda \tau)^{2} \tag{3.3}
\end{equation*}
$$

Differentianting equation (3.2) with respect to $s$ and using Frenet frame (3.2), we obtain

$$
\begin{equation*}
\kappa^{*} N^{*}\left(f^{\prime}\right)^{2}+f^{\prime \prime} T^{*}=\left(-\lambda \kappa^{\prime}\right) T+\left(\kappa-\lambda \kappa^{2}-\lambda \tau^{2}\right) N+\left(\lambda \tau^{\prime}\right) B \tag{3.4}
\end{equation*}
$$

By taking the inner product of the last relation with $N=\cos \theta N^{*}+\sin \theta B^{*}$, we get

$$
\begin{equation*}
\kappa^{*}\left(f^{\prime}\right)^{2} \cos \theta=\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right) . \tag{3.5}
\end{equation*}
$$

Then, by using equation (3.3), we find

$$
\begin{equation*}
\kappa^{*}=\frac{\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)}{\left[(1-\lambda \kappa)^{2}+\lambda^{2} \tau^{2}\right] \cos \theta} \tag{3.6}
\end{equation*}
$$

Putting the equations (3.2) and (3.6) in (3.4), we get
where, $e=\frac{-\cos \theta\left(\lambda \kappa^{\prime}+\frac{f^{\prime \prime}}{f^{\prime}}(1-\lambda \kappa)\right)}{\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)} \quad$ and $\quad h=\frac{\cos \theta\left(\lambda \tau^{\prime}+\frac{f^{\prime \prime}}{f^{\prime}} \lambda \tau\right)}{\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)}$.
Differentianting (3.7) with respect to $s$ and using Frenet frame (2.1), we obtain

$$
\begin{equation*}
\left(-\kappa^{*} T^{*}+\tau^{*} B^{*}\right) f^{\prime}=\left(e^{\prime}-\kappa \cos \theta\right) T+(e \kappa-h \tau) N+\left(\tau \cos \theta+h^{\prime}\right) B \tag{3.8}
\end{equation*}
$$

By using (3.2) and (3.8), we get

$$
\begin{equation*}
\tau^{*} B^{*}=\frac{D_{1}}{f^{\prime}} T+\frac{D_{2}}{f^{\prime}} N+\frac{D_{3}}{f^{\prime}} B \tag{3.9}
\end{equation*}
$$

where

$$
\begin{aligned}
& D_{1}=e^{\prime}-\kappa \cos \theta+\kappa^{*}(1-\lambda \kappa) \\
& D_{2}=e \kappa-h \tau \\
& D_{3}=\tau \cos \theta+h^{\prime}+\kappa^{*} \lambda \tau
\end{aligned}
$$

By taking the inner product of equation (3.9) with itself, we get

$$
\begin{equation*}
T^{*}=\frac{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}}{f^{\prime}} \tag{3.10}
\end{equation*}
$$

By using equation (3.10) in equation (3.9), we obtain

$$
\begin{equation*}
B^{*}=\frac{D_{1}}{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}} T+\frac{D_{2}}{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}} N+\frac{D_{3}}{\sqrt{D_{1}^{2}+D_{2}^{2}+D_{3}^{2}}} B \tag{3.11}
\end{equation*}
$$

This completes the proof.

Theorem 3.2. Let $\alpha$ be a unit speed curve with $N$ principal normal vector in $\mathbb{E}^{3}$. $\alpha^{*}$ be a regular curve with $N^{*}$ principal normal vector, then $\left(\alpha, \alpha^{*}\right)$ is a pair of generalized Bertrand curve if and only if the curvature $\kappa(s)$ and torsion $\tau(s)$ of $\alpha(s)$ satisfy;
$\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)=\cos \theta \frac{\left\{\left[\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)\right]^{2}\left[\lambda^{2} \tau^{2}+(1-\lambda \kappa)^{2}\right]+\left[\lambda \tau^{\prime}-\lambda^{2}\left(\kappa \tau^{\prime}-\kappa^{\prime} \tau\right)\right]^{2}\right\}^{\frac{1}{2}}}{\left[(1-\lambda \kappa)^{2}+\lambda^{2} \tau^{2}\right]^{\frac{1}{2}}}$,
where $\theta$ is the angle between the vectors $N(s), N^{*}\left(s^{*}\right)$ and $\lambda$ is a non-zero constant,$\kappa^{\prime}=\frac{d \kappa}{d s}$ and $\tau^{\prime}=\frac{d \tau}{d s}$.

Proof. We assume that $\left(\alpha, \alpha^{*}\right)$ is a pair of generalized Bertrand curve in $\mathbb{E}^{3}$, then we have

$$
\begin{equation*}
\alpha^{*}\left(s^{*}\right)=\alpha(s)+\lambda(s) N(s) . \tag{3.13}
\end{equation*}
$$

From Theorem 1, we get

$$
\begin{equation*}
T^{*} f^{\prime}=(1-\lambda \kappa) T+\lambda \tau B \tag{3.14}
\end{equation*}
$$

and we have known,

$$
\begin{equation*}
N=\cos \theta N^{*}+\sin \theta B^{*} \tag{3.15}
\end{equation*}
$$

Differentiating (3.14) with respect to $s$ and using Frenet frame (2.1), we find

$$
\begin{equation*}
-\kappa T+\tau B=-\kappa^{*} \cos \theta T^{*} f^{\prime}-\tau^{*} \sin \theta N^{*} f^{\prime}+\tau^{*} \cos \theta B^{*} f^{\prime} \tag{3.16}
\end{equation*}
$$

By inner product (3.16) with (3.14), we reached

$$
\begin{equation*}
\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)=\kappa^{*} \cos \theta\left(f^{\prime}\right)^{2} \tag{3.17}
\end{equation*}
$$

The curvature $\kappa^{*}\left(s^{*}\right)$ of the curve $\alpha^{*}\left(s^{*}\right)$ is

$$
\kappa^{*}\left(s^{*}\right)=\frac{\left\|\alpha^{*^{\prime}} \times \alpha^{*^{\prime \prime}}\right\|}{\left\|\alpha^{*^{\prime}}\right\|^{3}}
$$

where, $\alpha^{*^{\prime}}=\frac{d \alpha^{*}}{d s^{*}} \frac{d s^{*}}{d s}=\frac{d \alpha^{*}}{d s},\left\|\alpha^{*^{\prime}}\right\|=\left\|\frac{d \alpha^{*}}{d s}\right\|=\left\|f^{\prime}\right\|$.
By using the last relations in (3.17), we easily get

$$
\begin{equation*}
\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)=\frac{\left\|\alpha^{*^{\prime}} \times \alpha^{*^{\prime \prime}}\right\|}{\left\|\alpha^{*^{\prime}}\right\|} \cos \theta \tag{3.18}
\end{equation*}
$$

Also we have

$$
\alpha^{*^{\prime}}=(1-\lambda \kappa) T+\lambda \tau B
$$

and differentiating the last relation with respect to $s$, we get

$$
\alpha^{*^{\prime \prime}}=\lambda \kappa^{\prime} T+\left[\kappa-\left(\kappa^{2}+\tau^{2}\right)\right] N+\lambda \tau^{\prime} B
$$

Then if we calculate $\left\|\alpha^{*^{\prime}}\right\|$ and $\left\|\alpha^{*^{\prime}} \times \alpha^{*^{\prime \prime}}\right\|$, we get

$$
\begin{align*}
\left\|\alpha^{\alpha^{\prime}}\right\|= & {\left[(1-\lambda \kappa)^{2}+\lambda^{2} \tau^{2}\right]^{\frac{1}{2}}, } \\
\left\|\alpha^{*^{\prime}} \times \alpha^{*^{\prime \prime}}\right\|= & \left\{\lambda^{2} \tau^{2}\left[\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)\right]^{2}+\left[\lambda \tau^{\prime}-\lambda^{2}\left(\kappa \tau^{\prime}-\kappa^{\prime} \tau\right)\right]^{2}+\right.  \tag{3.19}\\
& {\left.\left[(1-\lambda \kappa)^{2}\left[\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)\right]^{2}\right]\right\}^{\frac{1}{2}} }
\end{align*}
$$

We put equations (3.19) in (3.18), we get equation (3.12) .
Conversely, we will prove that if $\kappa(s)$ and $\tau(s)$ satisfy equation (3.12), the principal normal and binormal of $\alpha^{*}$ generated by the equation

$$
\begin{equation*}
\alpha^{*}\left(s^{*}\right)=\alpha(s)+\lambda(s) N(s) \tag{3.20}
\end{equation*}
$$

are coplanar with the principal normal of $\alpha(s)$, where $s^{*}=s^{*}(s)$. The angle between $N$ and $N^{*}$ is $\theta$ in equation (3.12), we have known that $\lambda$ is a non-zero constant. Then from (3.20) differentiating with respect to $s$ we easily get

$$
\begin{equation*}
T^{*} f^{\prime}=(1-\lambda \kappa) T+\lambda \tau B \tag{3.21}
\end{equation*}
$$

By taking the inner product of equation (3.21) with $N$, we get

$$
\left\langle T^{*}, N\right\rangle=0
$$

which means that; $N$ is coplanar with $N^{*}$ and $B^{*}$.
Then we prove that

$$
\left\langle N, N^{*}\right\rangle=\cos \theta
$$

We assume that

$$
\begin{equation*}
N=a N^{*}+b B^{*} \quad a, b \in \mathbb{R} \tag{3.22}
\end{equation*}
$$

Differentiating equation (3.22) with respect to $s$, we find

$$
\begin{equation*}
-\kappa T+\tau B=-\kappa^{*} a T^{*} f^{\prime}-\tau^{*} b N^{*} f^{\prime}+\tau^{*} a B^{*} f^{\prime} \tag{3.23}
\end{equation*}
$$

By taking product equation (3.23) with equation (3.21), we get

$$
\begin{equation*}
\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)=\kappa^{*} a\left(f^{\prime}\right)^{2} . \tag{3.24}
\end{equation*}
$$

From equation (3.24) and (3.17) we easily obtain $\cos \theta=a$.
From equation (3.22) we get

$$
\left\langle N, N^{*}\right\rangle=a=\cos \theta
$$

This completes the proof.
Remark 3.1. In Theorem 2, when $\theta=0$, we have $\cos \theta=1$. Then we have
$\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)=\frac{\left\{\left[\kappa-\lambda\left(\kappa^{2}+\tau^{2}\right)\right]^{2}\left[\lambda^{2} \tau^{2}+(1-\lambda \kappa)^{2}\right]+\left[\lambda \tau^{\prime}-\lambda^{2}\left(\kappa \tau^{\prime}-\kappa^{\prime} \tau\right)\right]^{2}\right\}^{\frac{1}{2}}}{\left[(1-\lambda \kappa)^{2}+\lambda^{2} \tau^{2}\right]^{\frac{1}{2}}}$.
Squaring both sides of this equation, we find

$$
\left[\lambda \tau^{\prime}-\lambda^{2}\left(\kappa \tau^{\prime}-\kappa^{\prime} \tau\right)\right]=0
$$

Therefore ,

$$
\begin{aligned}
\tau^{\prime}(\lambda \kappa-1) & =\lambda \kappa^{\prime} \tau \\
\frac{d \tau}{\tau} & =\frac{d(\lambda \kappa-1)}{(\lambda \kappa-1)} \\
\lambda_{1} \tau+\lambda_{2} \kappa & =1
\end{aligned}
$$

where $\lambda$ is a constant. The equation $\lambda_{1} \tau+\lambda_{2} \kappa=1$ is the necessary and sufficient condition for a curve to be a Bertrand curve in Euclidean 3-space.

Example 3.1. Let us take $\theta=\frac{\pi}{4}, \lambda=2$ and $\kappa=1$ in Theorem 2, then we obtain $\tau=\frac{3 \tanh \left(\sqrt{\frac{5}{2}} s\right)}{\sqrt{10-4 \tanh \left(\sqrt{\frac{5}{2}} s\right)^{2}}}$. Thus we have generalized Bertrand curve $\alpha$ with curvatures $\kappa=1$ and $\tau=\frac{3 \tanh \left(\sqrt{\frac{5}{2}} s\right)}{\sqrt{10-4 \tanh \left(\sqrt{\frac{5}{2}} s\right)^{2}}}$ in Euclidean 3-space. It is easily check that $\alpha$ is a Salkowski curve.

Remark 3.2. If we take $\theta=\frac{\pi}{2}$ in equation 3.15, we get $N=B^{*}$ which means that $\alpha$ is a Mannheim curve and $\left(\alpha, \alpha^{*}\right)$ is a pair Mannheim curve. Also from Theorem 2, we get $\kappa=\lambda\left(\kappa^{2}+\tau^{2}\right)$. This equation also shows that $\alpha$ is a Mannheim curve .

Remark 3.3. We know that curves with constant curvatures (circular helix) are Bertrand curves. The equation (3.12) satisfies the circular helix if and only if $\theta=0$ or $\theta=\pi$ .In this case, circular helices are not generalized by Bertrand curves but only by classical Bertrand curves.

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