# FINSLER SPACE SUBJECTED TO A KROPINA CHANGE WITH AN $h$-VECTOR 

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#### Abstract

In this paper, we discuss the Finsler spaces $\left(M^{n}, L\right)$ and $\left(M^{n},{ }^{*} L\right)$, where ${ }^{*} L(x, y)$ is obtained from $L(x, y)$ by Kropina change ${ }^{*} L(x, y)=\frac{L^{2}(x, y)}{b_{i}(x, y) y^{i}}$ and $b_{i}(x, y)$ is an $h$-vector in $\left(M^{n}, L\right)$. We find the necessary and sufficient condition when the Cartan connection coefficients for both spaces $\left(M^{n}, L\right)$ and $\left(M^{n},{ }^{*} L\right)$ are the same. We also find the necessary and sufficient condition for Kropina change with an $h$-vector to be projective. Keywords: Finsler space, Kropina change, $h$-vector.


## 1. Introduction

In 1984, C. Shibata [16] dealt with a change of Finsler metric which is called a $\beta$ change of metric. A remarkable class of $\beta$-change is Kropina change $L(x, y)=\frac{L^{2}(x, y)}{b_{i}(x) y^{y}}$. If $L(x, y)$ is a metric function of a Riemannian space then ${ }^{*} L(x, y)$ reduces to the metric function of a Kropina space. Kropina metric was first introduced by L. Berwald in connection with a two-dimensional Finsler space with rectilinear extremals and was investigated by V.K. Kropina [8, 9]. Kropina metric is the simplest non-trivial Finsler metric having many interesting applications in physics, electron optics with a magnetic field, plants study for fungal fusion hypothesis, dissipative mechanics and irreversible thermodynamics [1, 2, 3, 6]. In 1978, C. Shibata [15] studied some basic local geometric properties of Kropina spaces. In 1991, M. Matsumoto obtained a set of necessary and sufficient conditions for a Kropina space to be of constant curvature [12].
H. Izumi [7], while studying the conformal transformation of Finsler spaces, introduced the concept of $h$-vector $b_{i}$, which is $v$-covariant constant with respect to the Cartan connection and satisfies $L C_{i j}^{h} b_{h}=\rho h_{i j}$, where $\rho$ is a non-zero scalar function, $C_{i j}^{h}$ are components of Cartan tensor and $h_{i j}$ are components of angular metric tensor. Thus if $b_{i}$ is an $h$-vector then
(i) $\left.b_{i}\right|_{k}=0$,
(ii) $L C_{i j}^{h} b_{h}=\rho h_{i j}$.

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This gives

$$
\begin{equation*}
L \dot{\partial}_{j} b_{i}=\rho h_{i j} \tag{1.2}
\end{equation*}
$$

Since $\rho \neq 0$ and $h_{i j} \neq 0$, the $h$-vector $b_{i}$ depends not only on positional coordinates but also on directional arguments. Izumi [7] showed that $\rho$ is independent of directional arguments. M. Matsumoto [11] discussed the Cartan connection of Randers change of Finsler metric, while B. N. Prasad [14] obtained the Cartan connection of $\left(M^{n},{ }^{*} L\right)$ where ${ }^{*} L(x, y)$ is given by ${ }^{*} L(x, y)=L(x, y)+b_{i}(x, y) y^{i}$, and $b_{i}(x, y)$ is an $h$-vector. Present authors [4,5] discussed the hypersurface of a Finsler space whose metric is given by certain transformations with an $h$-vector. In this paper we obtain the relation between the Cartan connections of $F^{n}=\left(M^{n}, L\right)$ and ${ }^{*} F^{n}=\left(M^{n},{ }^{*} L\right)$ where ${ }^{*} L(x, y)$ is obtained by the transformation

$$
\begin{equation*}
{ }^{*} L(x, y)=\frac{L^{2}(x, y)}{b_{i}(x, y) y^{i}} \tag{1.3}
\end{equation*}
$$

and $b_{i}(x, y)$ is an $h$-vector in $\left(M^{n}, L\right)$.
The paper is organized as follows: In Section 2, we study how the fundamental metric tensor and the Cartan tensor change by Kropina change with an $h$-vector. The relation between the Cartan connection coefficients of both spaces is obtained in the Section 3 and we find the necessary and sufficient condition when these connection coefficients are the same. In Section 4, we find the necessary and sufficient condition for the Kropina change with an $h$-vector to be projective.

## 2. The Finsler space ${ }^{*} F^{n}=\left(M^{n},{ }^{*} L\right)$

Let $F^{n}=\left(M^{n}, L\right)$ be an $n$-dimensional Finsler space equipped with the fundamental function $L(x, y)$. We consider a change of Finsler metric ${ }^{*} L(x, y)$ which is defined by (1.3) and have another Finsler space ${ }^{*} F^{n}=\left(M^{n},{ }^{*} L\right)$. If we denote $b_{i} y^{i}$ by $\beta$ then the indicatory property of $h_{i j}$ yields $\dot{\partial}_{i} \beta=b_{i}$. Throughout this paper, the geometric objects associated with ${ }^{*} F^{n}$ will be marked by the asterisk. We shall use the notation

$$
L_{i}=\dot{\partial}_{i} L=l_{i}, L_{i j}=\dot{\partial}_{i} \dot{\partial}_{j} L, L_{i j k}=\dot{\partial}_{k} L_{i j}, \ldots
$$

etc. From (1.3), we get

$$
\begin{gather*}
{ }^{*} L_{i}=2 \tau L_{i}-\tau^{2} b_{i},  \tag{2.1}\\
{ }^{*} L_{i j}=\left(2 \tau-\rho \tau^{2}\right) L_{i j}+\frac{2 \tau^{2}}{\beta} m_{i} m_{j},  \tag{2.2}\\
{ }^{*} L_{i j k}=\left(2 \tau-\rho \tau^{2}\right) L_{i j k}+\frac{2 \tau}{\beta}(\rho \tau-1)\left(m_{i} L_{j k}+m_{j} L_{i k}+m_{k} L_{i j}\right) \\
-\frac{2 \tau^{2}}{L \beta}\left(m_{i} m_{j} l_{k}+m_{j} m_{k} l_{i}+m_{k} m_{i} l_{j}\right)-\frac{6 \tau^{2}}{\beta^{2}} m_{i} m_{j} m_{k}, \tag{2.3}
\end{gather*}
$$

where $\tau=L / \beta, m_{i}=b_{i}-\frac{1}{\tau} l_{i}$. The normalized supporting element, the metric tensor and Cartan tensor of ${ }^{*} F^{n}$ are obtained as

$$
\begin{gather*}
{ }^{*} l_{i}=2 \tau l_{i}-\tau^{2} b_{i},  \tag{2.4}\\
{ }^{*} g_{i j}=\left(2 \tau^{2}-\rho \tau^{3}\right) g_{i j}+3 \tau^{4} b_{i} b_{j}-4 \tau^{3}\left(l_{i} b_{j}+b_{i} l_{j}\right)+\left(4 \tau^{2}+\rho \tau^{3}\right) l_{i} l_{j},  \tag{2.5}\\
{ }^{*} C_{i j k}=\left(2 \tau^{2}-\rho \tau^{3}\right) C_{i j k}-\frac{\tau^{2}}{2 \beta}(4-3 \rho \tau)\left(h_{i j} m_{k}+h_{j k} m_{i}+h_{k i} m_{j}\right)-\frac{6 \tau^{2}}{\beta} m_{i} m_{j} m_{k} . \tag{2.6}
\end{gather*}
$$

For the computation of the inverse metric tensor, we use the following lemma [10]:

Lemma 2.1. Let $\left(m_{i j}\right)$ be a non-singular matrix and $l_{i j}=m_{i j}+n_{i} n_{j ;}$. The elements $l^{i j}$ of the inverse matrix and the determinant of the matrix $\left(l_{i j}\right)$ are given by

$$
l^{i j}=m^{i j}-\left(1+n_{k} n^{k}\right)^{-1} n^{i} n^{j}, \quad \operatorname{det}\left(l_{i j}\right)=\left(1+n_{k} n^{k}\right) \operatorname{det}\left(m_{i j}\right)
$$

respectively, where $m^{i j}$ are elements of the inverse matrix of $\left(m_{i j}\right)$ and $n^{k}=m^{k i} n_{i}$.
The inverse metric tensor of ${ }^{\text {" }} F^{n}$ is derived as follows:

$$
{ }^{*} g^{i j}=\left(2 \tau^{2}-\rho \tau^{3}\right)^{-1}\left[g^{i j}-\frac{2 \tau}{2 b^{2} \tau-\rho} b^{i} b^{j}+\frac{4-\rho \tau}{2 b^{2} \tau-\rho}\left(l^{i} b^{j}+b^{i} l^{j}\right)\right.
$$

$$
\begin{equation*}
\left.-\frac{3 \rho b^{2} \tau^{3}-\rho^{2} \tau^{2}-4 b^{2} \tau^{2}-2 \rho \tau+8}{\tau\left(2 b^{2} \tau-\rho\right)} l^{i} l^{j}\right] \tag{2.7}
\end{equation*}
$$

where $b$ is the magnitude of the vector $b^{i}=g^{i j} b_{j}$.
From (2.6) and (2.7), we get

$$
\begin{align*}
{ }^{*} C_{i j}^{h}=C_{i j}^{h}-\frac{(4-3 \rho \tau) \tau}{2 L(2-\rho \tau)}\left(h_{i j} m^{h}\right. & \left.+h_{j}^{h} m_{i}+h_{i}^{h} m_{j}\right)-\frac{6 \tau}{L(2-\rho \tau)} m_{i} m_{j} m^{h} \\
+\frac{2 \tau b^{h}-(4-\rho \tau) l^{h}}{L(2-\rho \tau)\left(2 b^{2} \tau-\rho\right)} & {\left[h_{i j}\left\{\frac{1}{2} m^{2} \tau(4-3 \rho \tau)-\rho(2-\rho \tau)\right\}\right.}  \tag{2.8}\\
& \left.+m_{i} m_{j}\left\{6 \tau m^{2}+\tau(4-3 \rho \tau)\right\}\right] .
\end{align*}
$$

## 3. Cartan connection of the space ${ }^{*} F^{n}$

Let $C^{*} \Gamma=\left({ }^{*} F_{j k}^{i}{ }^{,}{ }^{N} N_{j}^{i},{ }^{*} C_{j k}^{i}\right)$ be the Cartan connection for the space ${ }^{*} F^{n}=\left(M^{n},{ }^{*}\right)$. Since for a Cartan connection $L_{i \mid j}=0$, we obtain

$$
\begin{equation*}
\partial_{j} L_{i}=L_{i r} N_{j}^{r}+L_{r} F_{i j}^{r} . \tag{3.1}
\end{equation*}
$$

Differentiating (2.1) with respect to $x^{j}$, we get

$$
\begin{equation*}
\partial_{j}{ }^{*} L_{i}=2 \tau \partial_{j} L_{i}+2 L_{i} \partial_{j} \tau-\tau^{2} \partial_{j} b_{i}-2 \tau b_{i} \partial_{j} \tau . \tag{3.2}
\end{equation*}
$$

This equation may be written in tensorial form as

$$
\begin{align*}
{ }^{*} L_{i r}{ }^{*} N_{j}^{r}+{ }^{*} L_{r}{ }^{*}{ }^{*} F_{i j}^{r} & =2 \tau\left(L_{i r} N_{j}^{r}+L_{r} F_{i j}^{r}\right)+\frac{2 L_{i}}{\beta}\left(N_{j}^{r} L_{r}-\tau \beta_{j}-\tau N_{j}^{r} b_{r}\right)  \tag{3.3}\\
& -\tau^{2}\left(b_{i l j}+\rho L_{i r} N_{j}^{r}+b_{r} F_{i j}^{r}\right)-\frac{2 \tau}{\beta} b_{i}\left(N_{j}^{r} L_{r}-\tau \beta_{j}-\tau N_{j}^{r} b_{r}\right),
\end{align*}
$$

where $\beta_{j}=\beta_{\mid j}$. If we put

$$
\begin{equation*}
{ }^{*} F_{j k}^{i}=F_{j k}^{i}+D_{j k}^{i}, \tag{3.4}
\end{equation*}
$$

then in view of (2.2), equation (3.3) may be written as

$$
\begin{equation*}
\left(2 \tau l_{r}-\tau^{2} b_{r}\right) D_{i j}^{r}+\left\{\left(2 \tau-\rho \tau^{2}\right) L_{i r}+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\} D_{0 j}^{r}=\frac{2 \tau^{2}}{\beta} m_{i} \beta_{j}-\tau^{2} b_{i \mid j}, \tag{3.5}
\end{equation*}
$$

where the subscript ' 0 ' denote the contraction by $y^{i}$.
In order to find the difference tensor $D_{j k^{k}}^{i}$, we construct supplementary equations to (3.5). From (2.2), we obtain

$$
\begin{align*}
\partial_{k}{ }^{*} L_{i j}= & \left(2 \tau-\rho \tau^{2}\right) \partial_{k} L_{i j}+L_{i j}\left(2 \partial_{k} \tau-2 \rho \tau \partial_{k} \tau-\tau^{2} \partial_{k} \rho\right) \\
& +\frac{2 \tau^{2}}{\beta} m_{i} \partial_{k} m_{j}+\frac{2 \tau^{2}}{\beta} m_{j} \partial_{k} m_{i}+2 m_{i} m_{j} \frac{1}{\beta^{2}}\left(2 \tau \beta \partial_{k} \tau-\tau^{2} \partial_{k} \beta\right) . \tag{3.6}
\end{align*}
$$

From $L_{i j k}=0$, equation (3.6) is written in the form

$$
\begin{aligned}
{ }^{*} L_{i j r}{ }^{*} N_{k}^{r}+{ }^{*} L_{r j}{ }^{*} F_{i k}^{r}+{ }^{*}{ }_{L i r}{ }^{*} F_{j k}^{r}= & \left(2 \tau-\rho \tau^{2}\right)\left\{L_{i j r} N_{k}^{r}+L_{r j} F_{i k}^{r}+L_{i r} F_{j k}^{r}\right\} \\
& +L_{i j}\left\{\frac{2}{\beta}(1-\rho \tau)\left(-\tau \beta_{k}-\tau N_{k}^{r} m_{r}\right)-\tau^{2} \rho_{k}\right\} \\
& +\frac{2 \tau^{2}}{\beta} m_{i}\left\{F_{j k}^{r} m_{r}+N_{k}^{r} L_{j r}\left(\rho-\frac{1}{\tau}\right)-\frac{1}{\beta \tau} N_{k}^{r} l_{j} m_{r}\right\} \\
& +\frac{2 \tau^{2}}{\beta} m_{j}\left\{F_{i k}^{r} m_{r}+N_{k}^{r} L_{i r}\left(\rho-\frac{1}{\tau}\right)-\frac{1}{\beta \tau} N_{k}^{r} l_{i} m_{r}\right\} \\
& +\frac{2}{\beta^{2}} m_{i} m_{j}\left\{2 \tau\left(-\tau \beta_{k}-\tau N_{k}^{r} m_{r}\right)-\tau^{2}\left(\beta_{k}+N_{k}^{r} b_{r}\right)\right\},
\end{aligned}
$$

where $\rho_{k}=\rho_{\mid k}=\partial_{k} \rho$. In view of (2.2), (2.3) and (3.4), above equation is written as

$$
\begin{align*}
& \left(2 \tau-\rho \tau^{2}\right)\left\{L_{i j r} D_{0 k}^{r}+L_{r j} D_{i k}^{r}+L_{i r} D_{j k}^{r}\right\}+\frac{2 \tau^{2}}{\beta} m_{r}\left(m_{j} D_{i k}^{r}+m_{i} D_{j k}^{r}\right) \\
& -\frac{2 \tau^{2}}{\beta^{2}}\left(m_{i} m_{j} b_{r}+m_{j} m_{r} b_{i}+m_{r} m_{i} b_{j}\right) D_{0 k}^{r}+\frac{2 \tau}{\beta} \beta_{k} L_{i j}-\frac{2 \rho \tau^{2}}{\beta} \beta_{k} L_{i j}  \tag{3.7}\\
& +\frac{2 \tau^{2}}{\beta}\left(\rho-\frac{1}{\tau}\right)\left(L_{j r} m_{i}+L_{i r} m_{j}+L_{i j} m_{r}\right) D_{0 k}^{r}+\frac{6 \tau^{2}}{\beta^{2}} \beta_{k} m_{i} m_{j}+\tau^{2} \rho_{k} L_{i j}=0 .
\end{align*}
$$

Now we will prove:
Proposition 3.1. The difference tensor $D_{j k}^{i}$ is completely determined by the equations (3.5) and (3.7).

To prove this, first we will prove a lemma:
Lemma 3.1. The system of algebraic equations
(i) ${ }^{*} L_{i r} A^{r}=B_{i}$,
(ii) ${ }^{*} L_{r} A^{r}=B$,
has a unique solution $A^{r}$ for given $B$ and $B_{i}$.
Proof. It follows from (2.2) that $(i)$ is written in the form

$$
\begin{equation*}
\left\{\left(2 \tau-\rho \tau^{2}\right) \frac{1}{L}\left(g_{i r}-l_{i} l_{r}\right)+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\} A^{r}=B_{i} . \tag{3.8}
\end{equation*}
$$

Contracting by $b^{i}$, we get

$$
\left\{\left(2 \tau-\rho \tau^{2}\right) \frac{1}{L}\left(b_{r}-\frac{1}{\tau} l_{r}\right)+\frac{2 \tau^{2}}{\beta} m^{2} m_{r}\right\} A^{r}=B_{\beta^{\prime}}
$$

i.e.

$$
\begin{equation*}
m_{r} A^{r}=B_{\beta}\left(\frac{2 \tau^{2} b^{2}}{\beta}-\frac{\rho \tau^{2}}{L}\right)^{-1}, \tag{3.9}
\end{equation*}
$$

where the subscript $\beta$ denote the contraction by $b^{i}$, i.e. $B_{\beta}=B_{i} b^{i}$. Also from (2.1), equation (ii) is written in the form

$$
\left(2 \tau l_{r}-\tau^{2} b_{r}\right) A^{r}=B
$$

i.e.

$$
\begin{equation*}
\tau^{2} m_{r} A^{r}-\tau l_{r} A^{r}=-B . \tag{3.10}
\end{equation*}
$$

Using (3.9) in (3.10), we get

$$
l_{r} A^{r}=\tau^{-1} B+\tau B_{\beta}\left(\frac{2 \tau^{2} b^{2}}{\beta}-\frac{\rho \tau^{2}}{L}\right)^{-1}
$$

Then (3.8) is written as
$g_{\text {ir }} A^{r}=\frac{L}{2 \tau-\rho \tau^{2}} B_{i}+l_{i}\left\{\tau^{-1} B+\tau B_{\beta}\left(\frac{2 \tau^{2} b^{2}}{\beta}-\frac{\rho \tau^{2}}{L}\right)^{-1}\right\}-\frac{2 \tau^{2}}{2-\rho \tau} m_{i} B_{\beta}\left(\frac{2 \tau^{2} b^{2}}{\beta}-\frac{\rho \tau^{2}}{L}\right)^{-1}$.
This gives

$$
\begin{equation*}
A^{i}=\frac{L}{2 \tau-\rho \tau^{2}} B^{i}+l^{i}\left\{\tau^{-1} B+\tau B_{\beta}\left(\frac{2 \tau^{2} b^{2}}{\beta}-\frac{\rho \tau^{2}}{L}\right)^{-1}\right\}-\frac{2 \tau^{2}}{2-\rho \tau} m^{i} B_{\beta}\left(\frac{2 \tau^{2} b^{2}}{\beta}-\frac{\rho \tau^{2}}{L}\right)^{-1}, \tag{3.11}
\end{equation*}
$$

which is the concrete form of the solution $A^{i}$.
We are now in a position to prove the proposition.
Taking the symmetric and anti-symmetric parts of (3.5), we get

$$
\begin{align*}
& 2\left(2 \tau l_{r}-\tau^{2} b_{r}\right) D_{i j}^{r}+\left\{\left(2 \tau-\rho \tau^{2}\right) L_{i r}+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\} D_{0 j}^{r} \\
& \quad+\left\{\left(2 \tau-\rho \tau^{2}\right) L_{j r}+\frac{2 \tau^{2}}{\beta} m_{j} m_{r}\right\} D_{0 i}^{r}=\frac{2 \tau^{2}}{\beta}\left(m_{i} \beta_{j}+m_{j} \beta_{i}\right)-2 \tau^{2} E_{i j}, \tag{3.12}
\end{align*}
$$

and

$$
\begin{gather*}
\left\{\left(2 \tau-\rho \tau^{2}\right) L_{i r}+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\} D_{0 j}^{r}-\left\{\left(2 \tau-\rho \tau^{2}\right) L_{j r}+\frac{2 \tau^{2}}{\beta} m_{j} m_{r}\right\} D_{0 i}^{r}  \tag{3.13}\\
=\frac{2 \tau^{2}}{\beta}\left(m_{i} \beta_{j}-m_{j} \beta_{i}\right)-2 \tau^{2} F_{i j}
\end{gather*}
$$

where we put $2 E_{i j}=b_{i \mid j}+b_{j \mid i}$ and $2 F_{i j}=b_{i \mid j}-b_{j \mid i}$.
On the other hand, applying Christoffel process with respected to indices $i, j, k$ in equation (3.7), we get

$$
\begin{align*}
& \left(2 \tau-\rho \tau^{2}\right)\left\{L_{i j r} D_{0 k}^{r}+L_{j k r} D_{0 i}^{r}-L_{k i r} D_{0 j}^{r}\right\}+2 D_{i k}^{r}\left\{\left(2 \tau-\rho \tau^{2}\right) L_{r j}+\frac{2 \tau^{2}}{\beta} m_{r} m_{j}\right\}  \tag{3.14}\\
& -\frac{2 \tau}{\beta}\left\{\beta_{k}\left((\rho \tau-1) L_{i j}-\frac{3 \tau}{\beta} m_{i} m_{j}\right)+\beta_{i}\left((\rho \tau-1) L_{j k}-\frac{3 \tau}{\beta} m_{j} m_{k}\right)-\beta_{j}\left((\rho \tau-1) L_{k i}-\frac{3 \tau}{\beta} m_{k} m_{i}\right)\right\} \\
& +\frac{2 \tau}{\beta} D_{0 k}^{r} \Im_{(i j r)} m_{i}\left\{(\rho \tau-1) L_{j r}-\frac{\tau}{\beta} m_{j} b_{r}\right\}+\frac{2 \tau}{\beta} D_{0 i}^{r} \Theta_{(j k r)} m_{j}\left\{(\rho \tau-1) L_{k r}-\frac{\tau}{\beta} m_{k} b_{r}\right\} \\
& -\frac{2 \tau}{\beta} D_{0 j}^{r} \Theta_{(k i r)} m_{k}\left\{(\rho \tau-1) L_{i r}-\frac{\tau}{\beta} m_{i} b_{r}\right\}+\tau^{2}\left(\rho_{k} L_{i j}+\rho_{i} L_{j k}-\rho_{j} L_{k i}\right)=0,
\end{align*}
$$

where $\Im_{(i j k)}$ denote cyclic interchange of indices $i, j, k$ and summation. Contracting (3.12) and (3.13) by $y^{j}$, we get

$$
\begin{equation*}
\left(4 \tau l_{r}-2 \tau^{2} b_{r}\right) D_{0 i}^{r}+\left\{\left(2 \tau-\rho \tau^{2}\right) L_{i r}+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\} D_{00}^{r}=\frac{2 \tau^{2}}{\beta} \beta_{0} m_{i}-2 \tau^{2} E_{i 0}, \tag{3.15}
\end{equation*}
$$

and

$$
\left\{\left(2 \tau-\rho \tau^{2}\right) L_{i r}+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\} D_{00}^{r}=\frac{2 \tau^{2}}{\beta} \beta_{0} m_{i}-2 \tau^{2} F_{i 0}
$$

i.e.

$$
\begin{equation*}
{ }^{*} L_{i r} D_{00}^{r}=\frac{2 \tau^{2}}{\beta} \beta_{0} m_{i}-2 \tau^{2} F_{i 0}, \tag{3.16}
\end{equation*}
$$

which on contraction by $b^{i}$ gives

$$
m_{r} D_{00}^{r}=\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1},
$$

where $\beta_{0}=\beta_{j} y^{j}$. Similarly contraction of (3.14) by $y^{k}$ gives

$$
\begin{align*}
& \left.\left(2 \tau-\rho \tau^{2}\right)\left\{L_{i j r} D_{00}^{r}-L_{j r} r_{0 i}^{r}+L_{i r} D_{0 j}^{r}\right\}+2 D_{0 i}^{r}\left(2 \tau-\rho \tau^{2}\right) L_{r j}+\frac{2 \tau^{2}}{\beta} m_{r} m_{j}\right\} \\
& +\frac{2 \tau}{\beta} D_{00}^{r} \Im_{(i j r)} m_{i}\left\{(\rho \tau-1) L_{j r}-\frac{\tau}{\beta} m_{j} b_{r}\right\}-\frac{2 \tau^{2}}{\beta} D_{0 i}^{r} m_{j} m_{r}+\frac{2 \tau^{2}}{\beta} D_{0 j}^{r} m_{i} m_{r}  \tag{3.17}\\
& -\frac{2 \tau}{\beta} \beta_{0}\left((\rho \tau-1) L_{i j}-\frac{3 \tau}{\beta} m_{i} m_{j}\right)+\tau^{2} \rho_{0} L_{i j}=0,
\end{align*}
$$

Contraction of (3.15) by $y^{i}$ gives

$$
\left(2 \tau l_{r}-\tau^{2} b_{r}\right) D_{00}^{r}=-\tau^{2} E_{00}
$$

i.e.

$$
\begin{equation*}
{ }^{*} L_{r} D_{00}^{r}=E_{00} . \tag{3.18}
\end{equation*}
$$

We can apply Lemma 3.1 to equations (3.16) and (3.18) to obtain

$$
\begin{align*}
D_{00}^{i}= & =l^{i} \tau\left\{\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1}-E_{00}\right\} \\
& -\frac{2 \tau^{2}}{2-\rho \tau}\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1} m^{i}+\frac{2 L \tau}{2-\rho \tau}\left(\frac{1}{\beta} \beta_{0} m^{i}-F_{0}^{i}\right), \tag{3.19}
\end{align*}
$$

where $F_{0}^{i}=g^{i j} F_{j 0}$. Also note that

$$
E_{00}=E_{i j} y^{i} y^{j}=b_{i \mid j} y^{i} y^{j}=\left(b_{i} y^{i}\right)_{\mid j} y^{j}=\beta_{10}=\beta_{0} .
$$

Now adding (3.13) and (3.17), we obtain

$$
D_{0 j}^{r}\left\{\left(2 \tau-\rho \tau^{2}\right) L_{i r}+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\}=G_{i j}
$$

i.e.

$$
\begin{equation*}
{ }^{*} L_{i r} D_{0 j}^{r}=G_{i j} \tag{3.20}
\end{equation*}
$$

where we put

$$
\begin{align*}
G_{i j} & =\frac{\tau^{2}}{\beta}\left(m_{i} \beta_{j}-m_{j} \beta_{i}\right)-\tau^{2} F_{i j}-\frac{1}{2}\left(2 \tau-\rho \tau^{2}\right) L_{i j r} D_{00}^{r}-\frac{\tau^{2}}{2} \rho_{0} L_{i j} \\
& -\frac{\tau}{\beta} D_{00}^{r} \Im_{(i j r)} m_{i}\left\{(\rho \tau-1) L_{j r}-\frac{\tau}{\beta} m_{j} b_{r}\right\}+\frac{\tau}{\beta} \beta_{0}\left\{(\rho \tau-1) L_{i j}-\frac{3 \tau}{\beta} m_{i} m_{j}\right\} \tag{3.21}
\end{align*}
$$

The equation (3.15) is written in the form

$$
\left(2 \tau l_{r}-\tau^{2} b_{r}\right) D_{0 j}^{r}=G_{j}
$$

i.e.

$$
\begin{equation*}
{ }^{*} L_{r} D_{0 j}^{r}=-\frac{1}{\tau^{2}} G_{j} \tag{3.22}
\end{equation*}
$$

where

$$
G_{j}=\frac{\tau^{2}}{\beta} \beta_{0} m_{j}-\tau^{2} E_{j 0}-\left\{\frac{1}{2}\left(2 \tau-\rho \tau^{2}\right) L_{j r}+\frac{\tau^{2}}{\beta} m_{j} m_{r}\right\} D_{00}^{r}
$$

In view of (3.16), $G_{j}$ are written as

$$
\begin{equation*}
G_{j}=\tau^{2}\left(F_{j 0}-E_{j 0}\right) \tag{3.23}
\end{equation*}
$$

Applying Lemma 3.1 to equations (3.20) and (3.22) to obtain

$$
\begin{equation*}
D_{0 j}^{i}=l^{i}\left\{\frac{1}{\tau}\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1} G_{\beta j}+\frac{1}{\tau} G_{j}\right\}-\frac{2 m^{i}}{2-\rho \tau}\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1} G_{\beta j}+\frac{L}{2 \tau-\rho \tau^{2}} G_{j^{\prime}}^{i} \tag{3.24}
\end{equation*}
$$

where $G_{j}^{i}=g^{i k} G_{k j}$.
Finally we solve (3.12) and (3.14) for $D_{j k}^{i}$. These equations may be written as

$$
\begin{equation*}
{ }^{*} L_{r j} D_{i k}^{r}=H_{j i k} \tag{3.25}
\end{equation*}
$$

and

$$
\begin{equation*}
{ }^{*} L_{r} D_{i k}^{r}=H_{i k} \tag{3.26}
\end{equation*}
$$

where

$$
\begin{align*}
H_{j i k} & =\frac{\left(\rho \tau^{2}-2 \tau\right)}{2}\left\{L_{i j r} D_{0 k}^{r}+L_{j k r} D_{0 i}^{r}-L_{k i r} D_{0 j}^{r}\right\}  \tag{3.27}\\
& -\frac{\tau}{\beta} D_{0 k}^{r} \Im_{(i j r)} m_{i}\left\{(\rho \tau-1) L_{j r}-\frac{\tau}{\beta} m_{j} b_{r}\right\}-\frac{\tau}{\beta} D_{0 i}^{r} \Im_{(j k r)} m_{j}\left\{(\rho \tau-1) L_{k r}-\frac{\tau}{\beta} m_{k} b_{r}\right\} \\
& +\frac{\tau}{\beta} D_{0 j}^{r} \Theta_{(k i r)} m_{k}\left\{(\rho \tau-1) L_{i r}-\frac{\tau}{\beta} m_{i} b_{r}\right\}-\frac{\tau^{2}}{2}\left(\rho_{k} L_{i j}+\rho_{i} L_{j k}-\rho_{j} L_{k i}\right) \\
& +\frac{\tau}{\beta}\left\{\beta_{k}\left((\rho \tau-1) L_{i j}-\frac{3 \tau}{\beta} m_{i} m_{j}\right)+\beta_{i}\left((\rho \tau-1) L_{j k}-\frac{3 \tau}{\beta} m_{j} m_{k}\right)\right. \\
& \left.-\beta_{j}\left((\rho \tau-1) L_{k i}-\frac{3 \tau}{\beta} m_{k} m_{i}\right)\right\},
\end{align*}
$$

and

$$
\begin{equation*}
H_{i k}=\frac{\tau^{2}}{\beta}\left(m_{i} \beta_{k}+m_{k} \beta_{i}\right)-\tau^{2} E_{i k}-\frac{1}{2}\left(G_{i k}+G_{k i}\right) \tag{3.28}
\end{equation*}
$$

Again applying Lemma 3.1 to equations (3.25) and (3.26) to obtain

$$
\begin{equation*}
D_{i k}^{j}=l^{j}\left\{\frac{1}{\tau}\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1} H_{\beta i k}+\frac{1}{\tau} H_{i k}\right\}-\frac{2 m^{j}}{2-\rho \tau}\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1} H_{\beta i k}+\frac{L}{2 \tau-\rho \tau^{2}} H_{i k}^{j} \tag{3.29}
\end{equation*}
$$

where we put $H_{i k}^{j}=g^{j m} H_{m i k}$. This completes the Proposition 3.1.
We now propose a lemma:
Lemma 3.2. If the h-vector is gradient then the scalar $\rho$ is constant.
Proof. Taking $h$-covariant derivative of (1.2) and using $L_{\mid k}=0$ and $h_{i j k}=0$, we get

$$
L\left(\dot{\partial}_{j} b_{i}\right)_{\mid k}=\rho_{\mid k} h_{i j}
$$

Utilizing the commutation formula exhibited by

$$
\dot{\partial}_{k}\left(T_{j \mid h}^{i}\right)-\left(\dot{\partial}_{k} T_{j}^{i}\right)_{\mid h}=T_{j}^{r} \dot{\partial}_{k} F_{r h}^{i}-T_{r}^{i} \dot{\partial}_{k} F_{j h}^{r}-\left(\dot{\partial}_{r} T_{j}^{i}\right) C_{h k \mid 0}^{r} ;
$$

we get

$$
2 L \dot{\partial}_{j} F_{i k}=\rho_{\mid k} h_{i j}-\rho_{\mid i} h_{j k} .
$$

If $b_{i}$ is a gradient vector, i.e. $2 F_{i j}=b_{i \mid j}-b_{j \mid i}=0$. Then above equation becomes

$$
\rho_{\mid k} h_{i j}-\rho_{\mid i} h_{j k}=0
$$

which after contraction by $y^{k}$ gives $\rho_{\mid k} y^{k}=0$. Differentiating $\rho_{\mid k} y^{k}=0$ partially with respect to $y^{j}$, and using the commutation formula $\dot{\partial}_{j}\left(\rho_{\mid k}\right)-\left(\dot{\partial}_{j} \rho\right)_{\mid k}=-\left(\dot{\partial}_{r} \rho\right) C_{j k \mid 0}^{r}$ and the fact that $\rho$ is a function of position only, we get $\rho_{\mid j}=0$ and therefore $\partial_{j} \rho=0$. This completes the proof.

Now, we find the condition for which the Cartan connection coefficients for both spaces $F^{n}$ and ${ }^{*} F^{n}$ are the same, i.e. ${ }^{*} F_{j k}^{i}=F_{j k}^{i}$ then $D_{j k}^{i}=0$. Therefore (3.15) and (3.16) gives $E_{i 0}=F_{i 0}$. This will give

$$
\begin{equation*}
b_{0 \mid i}=0, \tag{3.30}
\end{equation*}
$$

i.e. $\beta_{\mid i}=0$. Differentiating $\beta_{\mid i}=0$ partially with respect to $y^{j}$, and using the commutation formula $\dot{\partial}_{j}\left(\beta_{\mid i}\right)-\left(\dot{\partial}_{j} \beta\right)_{\mid i}=-\left(\dot{\partial}_{r} \beta\right) C_{i j 0}^{r}$, we get

$$
\begin{equation*}
b_{j \mid i}=b_{r} C_{i j \mid 0}^{r} . \tag{3.31}
\end{equation*}
$$

This gives $F_{i j}=0$ and then in view of Lemma 3.2, $F_{i j}=0$ implies $\rho_{i}=\rho_{\mid i}=0$.
Taking h-covariant derivative of (1.1)(ii) and using $L_{k k}=0, \rho_{\mid k}=0$ and $h_{i j k}=0$, we get $\left(b_{r} C_{i j}^{r}\right)_{\mid k}=\left(\frac{\rho}{L} h_{i j}\right)_{\mid k}=0$. This gives

$$
b_{r \mid k} C_{i j}^{r}+b_{r} C_{i j \mid k}^{r}=0 .
$$

From (3.31), we get $b_{r \mid k}=b_{k \mid r}$, then above equation becomes

$$
b_{k \mid r} C_{i j}^{r}+b_{r} C_{i j \mid k}^{r}=0 .
$$

Contracting by $y^{k}$, we get $b_{0 \mid r} C_{i j}^{r}+b_{r} C_{i j 0}^{r}=0$. Using (3.30) and (3.31), this gives $b_{i \mid j}=0$, i.e. the $h$-vector $b_{i}$ is parallel with respect to the Cartan connection of $F^{n}$.

Conversely, if $b_{i \mid j}=0$ then we get $E_{i j}=0=F_{i j}$ and $\beta_{i}=\beta_{\mid i}=b_{j \mid i} y^{j}=0$. In view of Lemma 3.2, $F_{i j}=0$ implies $\rho_{i}=\rho_{\mid i}=0$. Therefore from (3.19) we get $D_{00}^{i}=0$ and then $G_{i j}=0$ and $G_{j}=0$. This gives $D_{0 j}^{i}=0$ and then $H_{j i k}=0$ and $H_{i k}=0$. Therefore (3.29) implies $D_{j k}^{i}=0$ and then ${ }^{*} F_{j k}^{i}=F_{j k}^{i}$. Thus, we have:

Theorem 3.1. For the Kropina change with an h-vector, the Cartan connection coefficients for both spaces $F^{n}$ and ${ }^{*} F^{n}$ are the same if and only if the $h$-vector $b_{i}$ is parallel with respect to the Cartan connection of $F^{n}$.

Transvecting (3.4) by $y^{j}$ and using $F_{j k}^{i} y^{j}=G_{k^{\prime}}^{i}$ we get

$$
\begin{equation*}
{ }^{*} G_{k}^{i}=G_{k}^{i}+D_{0 k}^{i} . \tag{3.32}
\end{equation*}
$$

Further transvecting (3.32) by $y^{k}$ and using $G_{k}^{i} y^{k}=2 G^{i}$, we get

$$
\begin{equation*}
2^{*} G^{i}=2 G^{i}+D_{00}^{i} . \tag{3.33}
\end{equation*}
$$

Differentiating (3.32) partially with respect to $y^{h}$ and using $\dot{\partial}_{h} G_{k}^{i}=G_{k h}^{i}$, we have

$$
\begin{equation*}
{ }^{*} G_{k h}^{i}=G_{k h}^{i}+\dot{\partial}_{h} D_{0 k}^{i}, \tag{3.34}
\end{equation*}
$$

where $G_{k h}^{i}$ are the Berwald connection coefficients.
Now, if the $h$-vector $b_{i}$ is parallel with respect to the Cartan connection of $F^{n}$, then by Theorem 3.1, the Cartan connection coefficients for both spaces $F^{n}$ and ${ }^{*} F^{n}$ are the same, therefore $D_{j k}^{i}=0$. Then from (3.34), we get ${ }^{*} G_{k h}^{i}=G_{k h}^{i}$.
Thus, we have:
Theorem 3.2. For the Kropina change with an h-vector, if the h-vector $b_{i}$ is parallel with respect to the Cartan connection of $F^{n}$. Then the Berwald connection coefficients for both the spaces $F^{n}$ and $F^{n}$ are the same.

## 4. Relation between Projective change and Kropina change with an $h$-vector

We consider two Finsler spaces $F^{n}=\left(M^{n}, L\right)$ and ${ }^{*} F^{n}=\left(M^{n},{ }^{*} L\right)$. If any geodesic on $F^{n}$ is also a geodesic on ${ }^{*} F^{n}$ and the inverse is true, the change $L \rightarrow{ }^{*} L$ of the metric is called projective. A geodesic on $F^{n}$ is given by

$$
\frac{d y^{i}}{d t}+2 G^{i}(x, y)=\tau y^{i} ; \quad \tau=\frac{d^{2} s / d t^{2}}{d s / d t}
$$

The change $L \rightarrow{ }^{*}$ L is a projective change if and only if there exists a scalar $P(x, y)$ which is positively homogeneous of degree one in $y^{i}$ and satisfies [13]

$$
{ }^{*} G^{i}(x, y)=G^{i}(x, y)+P(x, y) y^{i} .
$$

Now, we find condition for the Kropina change (1.3) with $h$-vector to be projective. From (3.33), it follows that the Kropina change with an $h$-vector is projective if and only if $D_{00}^{i}=2 P y^{i}$. Then from (3.19), we get

$$
\begin{align*}
2 P y^{i}= & =l^{i} \tau\left\{\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1}-E_{00}\right\} \\
& -\frac{2 \tau^{2}}{2-\rho \tau}\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1} m^{i}+\frac{2 L \tau}{2-\rho \tau}\left(\frac{1}{\beta} \beta_{0} m^{i}-F_{0}^{i}\right) . \tag{4.1}
\end{align*}
$$

Contracting (4.1) by $y_{i}$ and using $m^{i} y_{i}=0=F_{0}^{i} y_{i}$, we get

$$
\begin{align*}
2 P L^{2} & =\tau\left\{\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1}-E_{00}\right\} L, \quad \text { i.e., } \\
P & =\frac{\tau}{2 L}\left\{\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1}-E_{00}\right\} . \tag{4.2}
\end{align*}
$$

Putting the value of $P$ in (4.1), we get

$$
-\frac{2 \tau^{2}}{2-\rho \tau}\left(\frac{2}{\beta} \beta_{0} m^{2}-2 F_{\beta 0}\right)\left(\frac{2}{\beta} b^{2}-\frac{\rho}{L}\right)^{-1} m^{i}+\frac{2 L \tau}{2-\rho \tau}\left(\frac{1}{\beta} \beta_{0} m^{i}-F_{0}^{i}\right)=0,
$$

i.e.,

$$
F_{0}^{i}=\frac{\beta_{0}}{\beta} m^{i}-\frac{1}{\beta} m_{r} D_{00}^{r} m^{i}
$$

Transvecing by $g_{i j}$, we get

$$
\begin{equation*}
F_{i 0}=\frac{\beta_{0}}{\beta} m_{i}-\frac{1}{\beta} m_{r} D_{00}^{r} m_{i} \tag{4.3}
\end{equation*}
$$

Using (4.3) in (3.16), and referring $2 \tau-\rho \tau^{2} \neq 0$, we get $L_{i r} D_{00}^{r}=0$, which transvecting by $m^{i}$ and using $L_{i r} m^{i}=\frac{1}{L} m_{r}$, we get $m_{r} D_{00}^{r}=0$, and then (4.3) becomes

$$
\begin{equation*}
F_{i 0}=\frac{\beta_{0}}{\beta} m_{i} \tag{4.4}
\end{equation*}
$$

This equation (4.4) is a necessary condition for the Kropina change with an $h$-vector to be a projective change.

Conversely, if (4.4) satisfies, then (3.16) gives

$$
\left\{\left(2 \tau-\rho \tau^{2}\right) L_{i r}+\frac{2 \tau^{2}}{\beta} m_{i} m_{r}\right\} D_{00}^{r}=0
$$

Transvecting by $m^{i}$ and referring $\frac{\left(2 \tau-\rho \tau^{2}\right)}{L}+\frac{2 \tau^{2}}{\beta} m^{2} \neq 0$, we get $m_{r} D_{00}^{r}=0$ and then (3.19) gives $D_{00}^{i}=-E_{00} \tau l^{i}$. Therefore ${ }^{*} F^{n}$ is projective to $F^{n}$. Thus, we have:

Theorem 4.1. The Kropina change (1.3) with an h-vector is projective if and only if the condition (4.4) is satisfied.

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