

A NEW LUMPED ELEMENT BRIDGED-T ABSORPTIVE BAND- STOP FILTER

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Abstract. *Following a brief review of previous work on bandstop filters, the inadequacy of a recent work to obtain a perfect notch or perfect absorption at the notch frequency ω_0 is demonstrated. A simple and elegant alternative solution, based on purely analytical arguments, is then presented. The resulting network is shown to achieve perfect matching as well as perfect absorption at the notch frequency and has several other advantages. A comparison has also been made with the conventional bridged-T band-stop filter.*

Key words: *bandstop filter, bridged-T network, circuit design.*

1. INTRODUCTION

Bandstop filters are circuits which reject, to within a specified tolerance, a band of frequencies around a centre frequency at which there is complete rejection. Such filters are known by various names, such as band rejection filters, notch filters, null networks etc. and are required in many situations in communication and instrumentation. Bandstop filters have fascinated a large number of researchers, including the present author, who has written papers on the analysis [1-7], design [8] and its limitations [9], and analysis and applications of dual input techniques to such filters [10-12]. All these contributions relate to analog circuits. Bandstop filters are also required in digital signal processing, and the author and his students have done extensive work on digital notch filters, using both FIR and IIR techniques [13-21]. Of these, [21] is a review of FIR notch filter design, which appeared in this journal.

At low frequencies, passive RC networks are mostly used, except in situations where a selectivity, defined as (notch frequency)/(3 dB stop bandwidth), is required to be more than half. In the latter cases, either active RC filters or LC networks are to be used. For high frequencies, LC networks are easily designed and implemented. At microwave frequencies, distributed networks are preferred over lumped networks, although the latter

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have the advantage of occupying less space, and as is well known, space is a premium in microwave integrated circuits. Examples of lumped element microwave bandstop filters can be found in [22-29], while bandstop filters with distributed elements can be found in [30,31].

2. SCOPE AND ORGANIZATION OF THE PAPER

This paper is concerned with the design of a band-stop filter which achieves a perfect notch and perfect absorption at some frequency ω_0 . In this context, we first demonstrate, in Section 3, the inadequacy of a recent solution proposed by Chieh and Rowland [32], by network theoretic arguments. In the next Section, we present a new, simple and elegant alternative design, based on purely analytical arguments. The resulting network is shown to achieve perfect matching as well as perfect absorption at the notch frequency, and has several other advantages. A normalized design is discussed in Section 5, and the simulation results are presented. A comparison of the new design with the conventional bridged-T band-stop filter is made in Section 6. Finally, Section 7 gives the concluding comments.

3. CHIEH AND ROWLAND'S DESIGN

Chieh and Rowland [32] proposed the symmetrical network of Fig. 1 where

$$Z_1(j\omega) = 1/(j\omega C), \quad (1a)$$

$$Z_2(j\omega) = R_1 + j\omega L_1 + 1/(j\omega C_1) \quad (1b)$$

and

$$Z_3(j\omega) = R_2 + j\omega L_2 + 1/(j\omega C_2). \quad (1c)$$

and both Z_2 and Z_3 resonate at the same frequency ω_0 . For ready reference, we reproduce here the expressions for the z -parameters of the network and the scattering parameters, in slightly different forms:

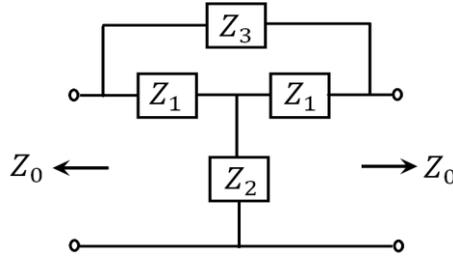


Fig. 1 The bridged-T network

$$z_{11} = z_{22} = Z_2 + (Z_1^2 + Z_1 Z_3) / (2Z_1 + Z_3), \quad (2)$$

$$z_{12} = z_{21} = Z_2 + Z_1^2 / (2Z_1 + Z_3), \quad (3)$$

$$S_{12} = S_{21} = 2z_{21} Z_0 / [(z_{11} + Z_0)^2 - z_{21}^2], \quad (4)$$

and

$$S_{11}=S_{22}=(z_{11}^2 - z_{21}^2)/[(z_{11}+Z_0)^2-z_{21}^2]. \quad (5)$$

Note from (2) and (3) that

$$z_{11}=z_{21}+Z_1Z_3/(2Z_1+Z_3). \quad (6)$$

From (4) and (5), we observe that for a perfect notch as well as perfect absorption at the frequency ω_0 , we require

$$z_{21}(j\omega_0)=0 \quad (7)$$

and

$$z_{11}(j\omega_0)=Z_0. \quad (8)$$

From (1), we have

$$Z_1(j\omega_0)=1/(j\omega_0C), Z_2(j\omega_0)=R_2, \text{ and } Z_3(j\omega_0)=R_1. \quad (9)$$

Substituting these values in (3) gives, on simplification,

$$z_{21}(j\omega_0)=R_2+1/[j\omega_0C(2+j\omega_0CR_1)], \quad (10)$$

which cannot be made zero. Also, under this condition,

$$z_{11}(j\omega_0)=R_2+(1+j\omega_0CR_1)/[j\omega_0C(2+j\omega_0CR_1)], \quad (11)$$

which cannot be equal to Z_0 if the latter is purely resistive, which is usually the case. Equation (8) can be satisfied only if Z_0 is a complex series RC impedance. Thus the network of Fig. 1 with the element values given by (1) can achieve neither perfect notch nor perfect absorption.

4. THE NEW DESIGN

The problem to be solved can be restated as follows: Given ω_0 and R_0 and the network topology of Fig.1, find Z_1 , Z_2 and Z_3 such that

$$z_{21}(j\omega_0)=Z_2(j\omega_0)+[Z_1(j\omega_0)]^2/[2Z_1(j\omega_0)+Z_3(j\omega_0)]=0, \quad (12)$$

and

$$z_{11}(j\omega_0)=z_{21}(j\omega_0)+Z_1(j\omega_0)Z_3(j\omega_0)/[2Z_1(j\omega_0)+Z_3(j\omega_0)]=R_0. \quad (13)$$

where Z_0 has been assumed to be resistive, equal to R_0 . In view of (12), (13) reduces to

$$z_{11}(j\omega_0)=Z_1(j\omega_0)Z_3(j\omega_0)/[2Z_1(j\omega_0)+Z_3(j\omega_0)]=R_0. \quad (14)$$

From (14), Z_3 is expressed in terms of Z_1 as

$$Z_3(j\omega_0)=2R_0Z_1(j\omega_0)/[Z_1(j\omega_0)-R_0]. \quad (15)$$

Combining this with (12) and simplifying, we get

$$Z_2(j\omega_0)=[R_0-Z_1(j\omega_0)]/2. \quad (16)$$

We can now choose a Z_1 . If we take $Z_1(j\omega_0)=1/(j\omega_0C)$, as in [1], then (15) gives, on simplification,

$$Z_3(j\omega_0)=[2R_0/(1+\omega_0^2C^2R_0^2)]+j\omega_0[2CR_0^2/(1+\omega_0^2C^2R_0^2)] \quad (17)$$

which represents a series combination of an inductance L_3 and a resistance R_3 , where

$$R_3 = [2R_0 / (1 + \omega_0^2 C^2 R_0^2)] \text{ and } L_3 = 2CR_0^2 / (1 + \omega_0^2 C^2 R_0^2). \quad (18)$$

Similarly, (16) gives

$$Z_2(j\omega_0) = (R_0/2) + j\omega_0 / (2\omega_0^2 C), \quad (19)$$

which also represents a series combination of an inductance L_2 and a resistance R_2 , where

$$R_2 = (R_0/2) \text{ and } L_2 = 1 / (2\omega_0^2 C). \quad (20)$$

In theory, C can be chosen to have any value, but as we shall see, it will be most convenient to choose C from the expression for R_3 given in (18), which gives

$$C = [(2R_0/R_3) - 1]^{1/2} / (\omega_0 R_0) \quad (21)$$

Note that if we choose

$$C = 1 / (\omega_0 R_0), \quad (22)$$

then R_3 becomes equal to R_0 . Also, under this condition, (17) and (18) give

$$L_3 = R_0 / \omega_0 \text{ and } L_2 = R_0 / (2\omega_0). \quad (23)$$

This choice of C is advantageous because then Z_3 can be obtained by a series combination of Z_2 and Z_2 and there is no spread in the element values of the network. Also note that lossy inductors can be used with ease because their losses can be absorbed in their series resistances. Finally, the element values of the network are consolidated as

$$C = 1 / (\omega_0 R_0), L_3 = 2L_2 = R_0 / \omega_0 \text{ and } R_3 = 2R_2 = R_0. \quad (24)$$

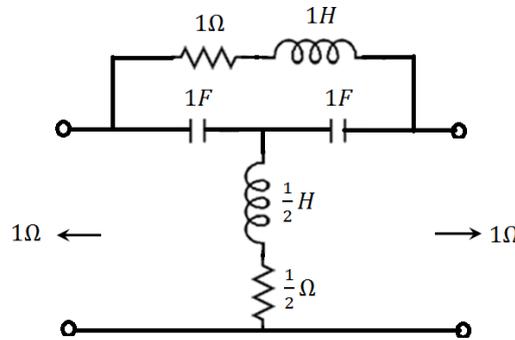


Fig. 2 The normalized design of the absorptive bandstop filter

5. A NORMALIZED DESIGN

It is always convenient to have a normalized design which can be denormalized by impedance and frequency scaling. Let $R_0 = 1$ ohm and $\omega_0 = 1$ rad/sec. Then (24) gives the element values as

$$C = 1F, L_3 = 2L_2 = 1H \text{ and } R_3 = 2R_2 = 1 \text{ ohm}. \quad (25)$$

The resulting network is shown in Fig. 2. This network has been simulated with MATLAB and the obtained plots of $|S_{11}(j\omega)|$ and $|S_{21}(j\omega)|$ are shown in Fig. 3. These plots exactly match the theoretical predictions.

6. COMPARISON WITH THE CONVENTIONAL BRIDGED-T BANDSTOP FILTER

It may be noted that compared to network proposed in [32], the conventional bridged-T bandstop filter [3] performs better because it achieves a perfect notch but not perfect absorption. In this network,

$$Z_1(j\omega_0)=1/(j\omega C), Z_2(j\omega_0)=r+j\omega L, \text{ and } Z_3(j\omega)=R. \quad (26)$$

The network then achieves a perfect notch at $\omega=[2/(LC)]^{1/2}$ under the condition $L=CRr$, but it cannot achieve $S_{11}(j\omega_0)=0$ unless Z_0 is a parallel combination of a capacitor C and a resistor $R/2$, which is not the usual case. Also, if we choose $r=R$, then there is no spread in the component values. Further, as in the proposed alternative, a lossy inductor can be used here. In addition, in comparison with the networks of [32] and that proposed here, it uses the least number, viz. three of reactive elements, yielding a transfer function of order three.

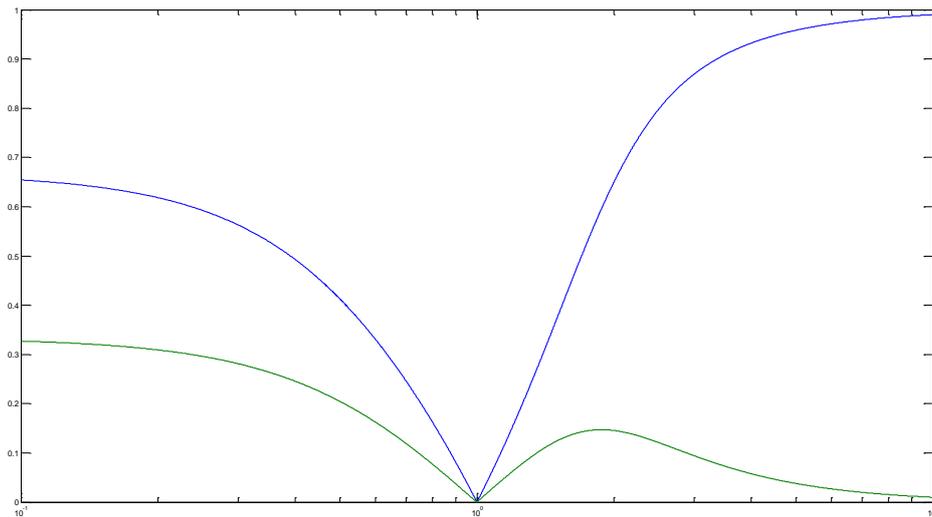


Fig. 3 Performance of the normalized design. The upper curve is a plot of $|S_{21}(j\omega)|$ and the lower curve represents $|S_{11}(j\omega)|$

7. CONCLUDING COMMENTS

It has been shown that the network proposed in [32] achieves neither a perfect notch nor perfect absorption. An alternative solution is proposed here purely by analytical, rather than physical or heuristic arguments, which achieves these two objectives simultaneously. The element values are obtained very simply, rather than by numerical and parametric methods as in [32]. Also, the new solution uses only two capacitors, instead of four, which

reduces the order of the transfer function by two. By an appropriate choice of the elements, there is no spread in the element values. A normalized design has been presented and the resulting characteristics of $|S_{11}(j\omega)|$ and $|S_{21}(j\omega_0)|$ have been plotted. A comparison of the two circuits has also been made with the conventional bridged-T bandstop filter.

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